

# 1 **Title**

2 Resolving the leftover problem for habitability and biosignatures  
3 on Enceladus

4 **Supplementary Material**

# 5 **Author list**

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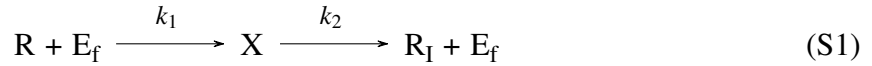
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# 1 Supplementary text

## 1.1 Mechanistic foundation for the Michaelis-Menten equation

The Michaelis-Menten (MM) approximation for uptake is typically obtained by considering the steady-state flow of a resource  $R$  (substrate) to the interior of the cell. This flow is decomposed into two steps, whereas the resource molecule first has to encounter a transporter protein located at the cellular surface, and then the transporter actively imports the molecule into the cytoplasm<sup>1-5</sup>:



where  $E_f$  is a ‘free’ porter,  $X$  is the porter-substrate complex, and  $R_I$  is the internal substrate.

While we illustrate here the derivation of the MM kinetics using the example of proteins transporting substrate to the cytoplasm (‘porters’), this also applies to substrates that are processed at the cellular surface without importation to the cytoplasm (e.g. possibly dihydrogen). Note that equation (S1) assumes that the unbinding of substrate and porter is negligible, and that importation (step 2) is irreversible. Using the law of mass action, assuming steady-state for  $X$  ( $\dot{X} = 0$ ) and setting  $[E] = [E_f] + [X]$ , an equivalent to familiar MM equation is obtained:

$$U(S) = \dot{R}_I = k_2[E] \frac{[R]}{[R] + \frac{k_2}{k_1}} \text{ mol s}^{-1} \quad (S2)$$

where  $k_2[E]$  ( $[E]$  counting the total number of mol of porter per cell) can be substituted for  $v_{max}$ , and  $k_2/k_1$  for  $K_m$ . However, equation (S2) does not entirely represent how the uptake rate is a function of the density of transporters on the cellular surface. Indeed,  $[E]$  also determines the rate of encounter between porters and substrate  $k_1$ <sup>6</sup>, such that the porter-specific rate of substrate capture  $k_1$  in units of  $\text{s}^{-1} (\text{mol L}^{-1})^{-1}$  can be approximated by:

$$k_1 = \rho_w \mathcal{N}_A 4\pi D_w r_c \frac{r_p}{r_p [E] \mathcal{N}_A + \pi r_c} \quad (S3)$$

where  $\mathcal{N}_A = 6.022 \times 10^{23} \text{ mol}^{-1}$  is Avogadro’s number,  $\rho_w = 1,000 \text{ L m}^{-3}$  converts  $\text{m}^3$  to L,  $D_w (\text{m}^2 \text{ s}^{-1})$  is the diffusion coefficient of  $R$  in water,  $r_c$  (m) is the cell’s radius (typically around  $1 \mu\text{m}$  for Archaea),  $r_p$  is the radius of the porter protein assumed to be disk-like (typically taken to be  $1 \text{ nm}^6$ ), and  $[E]$  (mol) is the cellular porter density.

## 62 1.2 Gradient analysis of $R^*$ under the MM-based model

63 Assuming Michaelis-Menten uptake, the limit resource value  $R^*$  follows (Box 1):

$$64 R^* = \frac{\gamma d K_m}{v_{max} - \gamma d} \quad (S4)$$

65 which is positive under the condition  $v_{max} - \gamma d$  (in other words that there exists a positive  
66 value of  $R$  such that growth occurs faster than mortality). We now turn to the sensitivity of  
67  $R^*$  to parameters assumed to change with temperature  $d$ ,  $K_m$ , and  $v_{max}$ . To this end, we look at  
68 the first order Taylor expansion of a small relative increase of each parameter taken separately  
69 using:

$$70 \xi_i = \frac{x_i}{R^*} \frac{\partial R^*}{\partial x_i} \quad (S5)$$

71 where  $x_i$  is a variable among  $d$ ,  $K_m$ , and  $v_{max}$ . By applying the expression in equation (S5) to  
72 equation (S4), we obtain:

$$73 \xi = \left( \xi_d = \frac{v_{max}}{v_{max} - \gamma d}, \xi_{K_m} = 1, \xi_{v_{max}} = -\frac{v_{max}}{v_{max} - \gamma d} \right). \quad (S6)$$

74 Under the condition that  $v_{max} > \gamma d > 0$ ,  $v_{max}/(v_{max} - \gamma d) > 1$ , implying that  $R^*$  is more  
75 sensitive to a relative increase in either  $v_{max}$  or  $d$  than in  $K_m$ . Interestingly,  $R^*$  is equally sensitive  
76 (with opposing signs) to a relative increase of  $d$  and to a relative decrease of  $v_{max}$ . Thus, if  
77 mortality and maximum uptake have the same sensitivity to temperature (for instance the  
78 same  $Q_{10}$ ),  $R^*$  would change according to the sensitivity of  $K_m$  to temperature, which may  
79 be either positive or negative in principle (see Main Text), although mechanistic principles  
80 underlying  $K_m$  indicate that it should increase with temperature (see Main Text), in accordance  
81 with observation<sup>7</sup>.

## 82 1.3 Activation energy of diffusion of $H_2$ in water

83 In order to approximate the activation energy for the diffusion constant of  $H_2$  in water, we  
84 used the temperature data compiled in the CRC Handbook for Chemistry and Physics, Table  
85 ‘Diffusion of Gases in Water’<sup>8</sup>. Using standard Arrhenius linear regression of  $\ln D_w \propto 1/T$ ,  
86 we obtained ( $R^2 > 0.99$ , two-sided Wald test  $p \approx 2 \times 10^{-8}$ ):

$$87 D_w^{H_2}(T) = (3.41 \times 10^{-6}) \times e^{-\frac{16,109}{RT}} \quad (S7)$$

88 i.e. that the diffusion constant of H<sub>2</sub> in water has an activation energy of 16.1 kJ mol<sup>-1</sup> (Supplementary  
89 Figure S1).

#### 90 **1.4 Activation energy of $v_{max}$**

91 Elias et al. (2014)<sup>9</sup> suggest that enzymes isolated from organisms adapted to various temperatures  
92 all have the same rate-temperature dependency, as captured by the empirical fold-increase of  
93 enzymatic rate per 10 K:

$$94 \quad Q_{10} = \frac{k_{cat}(T + 10)}{k_{cat}(T)} \quad (S8)$$

95 such that the temperature-rate scaling is approximated by:

$$96 \quad k_{cat}(T) = k_{cat}(T_{ref}) Q_{10}^{\frac{T-T_{ref}}{10}} \quad (S9)$$

97 The existence of a universal  $Q_{10}$  implies that the activation energies of enzymes vary with the  
98 temperature at which organisms are adapted via

$$99 \quad E_{a,c}(T_a) = RT_a \frac{T_a + 10}{10} \ln Q_{10} \quad (S10)$$

100 where  $T_a$  represents the temperature at which the organism from which the enzyme is isolated  
101 is adapted. This suggests that the activation energy of a Arrhenius function representing the  
102 rate-temperature scaling of  $v_{max}$  should range from 45 kJ mol<sup>-1</sup> at 273 K to 80 kJ mol<sup>-1</sup> at  
103 373 K, i.e. much higher than the activation energy for the diffusion constant for H<sub>2</sub> in water  
104 of 16 kJ mol<sup>-1</sup>.

#### 105 **1.5 Analysis of the relation between $R^*$ and $T$ in the yield-limit model**

106 Here, we analyze a generic version of the yield-limit model for any lithochemoautotrophic  
107 metabolism. Consider a generic catabolic reaction following the equation

$$108 \quad -e_D + \sum_i Y_i R_i = 0 \quad (S11)$$

109 where  $R_i$  are all the reagents (reactants and products) of the catabolic reaction with the exception  
110 of the electron donor  $e_D$  and  $Y_i$  are the stoichiometric terms (negative for reactants and positive  
111 for products) which are normalized for 1 mole of electron donor (such that the stoichiometric  
112 coefficient of  $e_D$  is  $-1$ ). Under the assumptions permitting to equate chemical activities with

113 concentrations, the Gibbs free energy per mole electron donor ( $\text{J mol}^{-1}$ ) for the catabolic reaction  
 114 is given by:

$$115 \quad \Delta_r G = \Delta_r G_0 + RT(\log \prod_i [R_i]^{Y_i} - \log[e_D]) = \Delta_r G_0 + RT \log Q' - RT \log[e_D] \quad (\text{S12})$$

116 where we have defined  $Q' := \prod_i [R_i]^{Y_i}$ , brackets denote the concentration ( $\text{mol L}^{-1}$ ),  $\Delta_r G_0$  is the  
 117 standard Gibbs free energy of reaction, which is usually known for a given reaction,  $R = 8.314$   
 118  $\text{JK}^{-1} \text{mol}^{-1}$  is the ideal gas constant and  $T$  is the temperature in Kelvin units.

119 Recalling the expression for the steady-state Gibbs free energy of reaction for the yield model  
 120 (Main Text equation (17)), neglecting the individual death rate unrelated to maintenance):

$$121 \quad (-\Delta_r G)^* = \frac{e_m(T)}{q_{cat}(T)} \quad (\text{S13})$$

122 we can solve for the electron donor concentration, yielding:

$$123 \quad [e_D]^*(T) = Q' \times \exp \left[ \frac{1}{RT} \left( \frac{e_m(T)}{q_{cat}(T)} + \Delta_r G_0 \right) \right] \quad (\text{S14})$$

124 where we recall that  $e_m$  is the biomass specific maintenance cost ( $\text{J mol}^{-1} \text{s}^{-1}$ ) and  $q_{cat}$  is the  
 125 biomass-specific catabolic rate, or rate of electron donor processing ( $\text{s}^{-1}$ ). Differentiating  
 126 equation (S14) with respect to temperature yields:

$$127 \quad \frac{([e_D]^*)'}{[e_D]^*} = \frac{1}{RT} \frac{e_m(T)}{q_{cat}(T)} \left( \frac{e'_m(T)}{e_m(T)} - \frac{q'_{cat}(T)}{q_{cat}(T)} - \frac{1}{T} \frac{\Delta_r G_0 q_{cat}(T)}{e_m(T)} \right). \quad (\text{S15})$$

128 Whether the threshold  $[e_D]^*$  increases with temperature is given by the sign of equation (S15).  
 129 For catabolic reactions able to power the metabolism, it is often the case that  $\Delta_r G_0 < 0$  (although  
 130 it is not strictly required for the catabolic reaction to power metabolism in physiological conditions),  
 131 in which case  $e'_m/e_m > q'_{cat}/q_{cat}$  guarantees positivity of equation (S15). The points brought  
 132 up in main text for  $e'_m/e_m > q'_{cat}/q_{cat}$  stand here as the molecular basis for these rates is the  
 133 same as  $v_{max}$  and  $d$  in the Michaelis-Menten based model. In the edge case that  $\Delta_r G_0$  is slightly  
 134 positive, the term in which it appears in equation (S15) should be vanishingly small, as the  
 135 decomposition of the energy rate  $e_m$  should give a product between an energy term much  
 136 greater than  $\Delta_r G_0$  ( $E_m$  is of the order of a  $\text{MJ mol}^{-1}$  for autotrophs) and a biomass-specific  
 137 rate term much larger than  $q_{cat}$ , owing to the scaling of the molecules which participate to  
 138 metabolism and those which contribute to maintenance cost (essentially the entire cell; see  
 139 ref.<sup>10</sup> for additional theoretical discussion). This ensures that even minute differences between

140  $e'_m/e_m$  and  $q'_{cat}/q_{cat}$  imply that  $[e_D]^*$  increases with temperature. These differences are likely  
 141 not minute for microorganisms living at their preferred temperature, as the asymmetrical characteristic  
 142 of the ‘thermal growth curve’ of microbes implies that the rate of maintenance/mortality increases  
 143 much faster than metabolic rates at the optimal growth temperature.

## 144 1.6 Uncertainty in P-H<sub>2</sub> co-limitation curves

145 Co-limitation curves (see Box 2) are defined by:

$$146 \frac{1}{\gamma_{H_2}} U_{H_2}([H_2]_{cl}) = \frac{1}{\gamma_P} U_P([P]_{cl}) \quad (S16)$$

147 where the subscript *cl* denotes concentrations at which co-limitation occurs. In our study, we  
 148 approximate uptake functions  $U_i$  using Michaelis-Menten kinetics:

$$149 U_i(R_i) = v_{max}^i \frac{R_i}{R_i + K_m^i} \quad (S17)$$

150 Thus, equation (S16) rewrites

$$151 [P]_{cl} = K_m^P \frac{\Gamma([H_2]_{cl})}{1 - \Gamma([H_2]_{cl})} \quad (S18)$$

152 where

$$153 \Gamma([H_2]) = \frac{\gamma_P v_{max}^{H_2}}{\gamma_{H_2} v_{max}^P} \frac{[H_2]}{[H_2] + K_m^{H_2}} \quad (S19)$$

154 Whether the product

$$155 \lim_{[H_2] \rightarrow \infty} \Gamma([H_2]) = \frac{\gamma_P v_{max}^{H_2}}{\gamma_{H_2} v_{max}^P} \quad (S20)$$

156 is greater or less than one qualitatively affects the shape of the co-limitation curve (equation  
 157 (S16)). In essence, this captures whether H<sub>2</sub> or P limits growth when both are abundant (i.e.  
 158 their concentration is much greater than their respective  $K_m$  values). Indeed, if equation (S20)  
 159 is lower than one, then  $[P]_{cl}$  should approach  $K_m^P(1 + \lim_{[H_2] \rightarrow \infty} \Gamma)$  when H<sub>2</sub> becomes much  
 160 greater than  $K_m^{H_2}$ , i.e growth tends to be limited by H<sub>2</sub> in ideal conditions. However, if  $\lim \Gamma >$   
 161 1, then there exists a value of  $[H_2]$  for which equation (S18) becomes indefinite ( $\Gamma = 1$ ), i.e.  
 162 growth tends to be P-limited in ideal conditions. In our three parameterizations, it is the latter  
 163 which is occurring, i.e. methanogens should be P-limited rather than H<sub>2</sub>-limited when both  
 164 are abundant. However, our empirical ranges for  $v_{max}^{H_2}$ ,  $v_{max}^P$ ,  $\gamma_{H_2}$  and  $\gamma_P$  could in principle  
 165 allow for combinations in which  $\Gamma_\infty > 1$ . Hence, that our three parameterizations have  $\Gamma_\infty$   
 166 rests on the particular selection of cases for the kinetic parameters for each resource, although

167 assuming that they would be correlated is not without merit (see Main text). In sum, we argue  
168 that we should typically expect hydrogenotrophic methanogens to be P rather than H<sub>2</sub> limited  
169 when both are abundant, but only empirical measurements of the uptake machinery for both  
170 resource in the same strain may resolve this. Speculatively, this should be expected to vary  
171 between strains and the environment from which they are isolated, as this behavior may be a  
172 by-product of adaptation of uptake machinery to conditions that typically differ widely from  
173 overabundance of every resource. As a note, such a difference between qualitative properties  
174 in the 2-resource system does not exist for the determination of the lines separating the basin  
175 in which trajectories reach the H<sub>2</sub> threshold (but not the P threshold) from the basin in which  
176 the opposite occurs.

## 177 **1.7 Monte-Carlo Markov Chain sampling of initial conditions that best** 178 **predict plume composition as final state**

179 **General approach** To identify concentrations which when used as initial conditions of our  
180 2-resource-1-population model (Main Text equations (11)) yield final, biotic, steady-states  
181 that best match the ranges of concentrations of H<sub>2</sub> and P inferred from chemical analysis of  
182 Enceladus's plume, we perform a Monte-Carlo Markov Chain (MCMC) sampling using Metropolis's  
183 algorithm. To do so, we first define the target distribution used to measure the likelihood of  
184 model outputs. This distribution is assumed to 2-dimensional log-normal, centered on the log-  
185 average of the ranges of assumed H<sub>2</sub> and P concentrations in the source material for the plume.  
186 In this case, we use the upper bound for the H<sub>2</sub> range ( $1 \times 10^{-2} \text{ mol kg}^{-1}$ ) as estimated in  
187 Glein and Truong (2025)<sup>11</sup>, and we take as lower bound the estimate by Waite et al. (2017)<sup>12</sup>  
188 ( $2 \times 10^{-7} \text{ mol kg}^{-1}$ ) for a log-average value of  $\approx 4.5 \times 10^{-5} \text{ mol kg}^{-1}$ . For P, our log-average  
189 is  $4.1 \times 10^{-3} \text{ mol kg}^{-1}$ . Then, we assume this multivariate log-normal distribution to have  
190 a diagonal covariance matrix where the diagonal elements are set such that the span (in log  
191 scale) of the ranges identified in Main Text Methods section *H<sub>2</sub> and P concentration ranges*  
192 captures 99% of the distribution's spread (as if the ranges corresponded to  $3\sigma$  error bars). This  
193 allows us to numerically evaluate the target probability density function, noted  $f_i$ , of proposed  
194 simulation final states ( $[\text{H}_2]_\infty, [\text{P}]_\infty$ ) using SciPy<sup>13</sup>.

195 **Model implementation** To simulate biotic steady-states as a function of boundary initial  
196 abiotic conditions, we first identify the minimal H<sub>2</sub> and P minimally viable concentration  
197 thresholds according to Main Text equation (9; Box 2) where the uptake function chosen to

198 be Michaelis-Menten:

$$\left\{ \begin{array}{l} [\text{H}_2]^* = \frac{\gamma_{\text{H}} d K_{\text{H}}}{v_{\text{max}}^{\text{H}} - \gamma_{\text{H}} d} \\ [\text{P}]^* = \frac{\gamma_{\text{P}} d K_{\text{P}}}{v_{\text{max}}^{\text{P}} - \gamma_{\text{P}} d} \end{array} \right. \quad \begin{array}{l} \text{(S21a)} \\ \text{(S21b)} \end{array}$$

200 using our smoothed approximation of parameter values (*Supplementary text*), where parameters  
201  $d$ ,  $\gamma_{\text{H}_2}$ ,  $\gamma_{\text{P}}$ ,  $K_m^{\text{H}_2}$ ,  $K_m^{\text{P}}$  are function of temperature (see *Construction of cases* in Main text Methods).

202 In doing so, we smoothly explore the range of resource thresholds between the minimum ( $[\text{H}_2]^* = 2.75 \times$   
203  $10^{-9} \text{ molL}^{-1}$  and  $[\text{P}]^* = 2.81 \times 10^{-11} \text{ molL}^{-1}$ ) and maximum ( $[\text{H}_2]^* = 5.6 \times 10^{-6} \text{ molL}^{-1}$  and  
204  $[\text{P}]^* = 2.71 \times 10^{-8} \text{ molL}^{-1}$ ) in our parameterizations (Supplementary Table S1).

205 Then, given proposed boundary abiotic initial conditions  $[\text{H}]_0$ ,  $[\text{P}]_0$  and temperature  $T$ , we first  
206 check whether they are viable ( $[\text{H}_2]_0 > [\text{H}_2]^*$  and  $[\text{P}]_0 > [\text{P}]^*$ ) and then identify the steady-state  
207 limiting resource which is  $\text{H}_2$  if

$$[\text{P}]_0 > [\text{H}_2]_0 \frac{\gamma_{\text{P}}}{\gamma_{\text{H}}} + [\text{P}]^* \quad \text{(S22)}$$

209 or P otherwise. Then, the steady-state concentration of the non-limiting resource is set by (see  
210 Box 2):

$$[\text{R}_n]_{\infty} = [\text{R}_n]_0 - \frac{\gamma_n}{\gamma_l} ([\text{R}_l]_0 - [\text{R}_l]^*) . \quad \text{(S23)}$$

212 This formally results in a numerical procedure to evaluate the function  $\pi : ([\text{H}_2]_0, [\text{P}]_0, T) \mapsto$   
213  $([\text{H}_2]_{\infty}, [\text{P}]_{\infty})$ .

214 **Monte Carlo Markov Chain sampling** We now use the model  $\pi$ , and the target probability  
215 density function  $f_t$  (defined by the plume observations) to feed into a Metropolis algorithm  
216 with a multivariate normal proposal distribution (log-scaled for  $[\text{H}_2]_0$ ,  $[\text{P}]_0$ , linear for  $T$ ) with  
217 standard deviations of 0.01, 0.01 and 3 respectively. We define our acceptance ratio as  $f_t(\pi(\mathbf{x}_p)) / f_t(\pi(\mathbf{x}_t))$   
218 where  $\mathbf{x}_t$  is the current state ( $[\text{H}_2]_0, [\text{P}]_0, T$ ) and  $\mathbf{x}_p$  the proposed state. We allow for a burn  
219 in period of 50,000 steps, then run the MCMC algorithm until 150,000 samples of initial  
220 conditions have been accepted.

221 **Plume best match hydrothermal  $\text{H}_2$  for hydrogenotrophic methanogenesis** Affholder et  
222 al. (2021)<sup>14</sup> describe a model of mixing of ocean water and hydrothermal fluid designed to  
223 simulate hydrothermal methanogens at Enceladus's seafloor. This model of mixing calculates  
224 the overall output of the hydrothermal environment as fluxes of  $\text{H}_2$  and  $\text{CH}_4$  being released to  
225 the ocean as a upward-rising plume. Affholder et al. (2022)<sup>14</sup> simulate values of these fluxes  
226 under various conditions of hydrothermal fluid and ocean temperatures and compositions

227 (concentrations of  $\text{CO}_2$ ,  $\text{CH}_4$ , and  $\text{H}_2$ ), as well as by including or not the feedback of hydrogenotrophic  
228 methanogens on their local environment of the mixing layer at the core-ocean interface. Then,  
229 under the assumption that the bulk concentration of  $\text{H}_2$  and  $\text{CH}_4$  in Enceladus's ocean are at  
230 steady-state, these fluxes are compared with the escape fluxes of  $\text{H}_2$  and  $\text{CH}_4$  in Enceladus's  
231 space plume as constrained by measurements carried out by the Cassini mission<sup>12</sup>. Affholder  
232 et al. (2022)<sup>15</sup> then construct a 'posterior' set of biotic simulations, by retaining the 500 simulations  
233 closest (by euclidean distance in the  $\text{H}_2$  and  $\text{CH}_4$  escape fluxes space) to Cassini observations<sup>12</sup>.  
234 By doing so, they identified the set of parameters (hydrothermal fluid and ocean temperature  
235 and concentrations of  $\text{H}_2$ ,  $\text{CO}_2$ , and  $\text{CH}_4$ ) that yield the simulations that best match plume  
236 observables when methanogens are included in the model. By fitting a gaussian curve to the  
237 values of  $\text{H}_2$  concentrations in the hydrothermal fluid used in the simulations that make up  
238 the 'posterior' set, a 95% interval of the 'posterior'  $\text{H}_2$  concentration in the hydrothermal  
239 fluid is obtained (section 2.5 of ref.<sup>15</sup>). This interval is 100-600  $\text{mmolL}^{-1}$ . Of important  
240 note, this concentration is that of the hydrothermal fluid used in the simulations of ref.<sup>14</sup>,  
241 however, in those simulations methanogens may locally experience any concentration between  
242 that of the hydrothermal fluid and that of the ocean, as they simulate a radial gradient. The  
243 ocean concentration in ref.<sup>14</sup> is taken to be lower than that of the hydrothermal fluid, thus,  
244 the 'posterior' range used here corresponds to the higher upper limit of the concentration that  
245 methanogens would experience locally at the seafloor, and does not correspond to that of the  
246 bulk of the ocean.

## 247 **1.8 Global feedback of hydrothermal, energy-limited, methanogens on P** 248 **budget in the bulk ocean**

249 An upper limit to biomass productivity of energy-limited hypothetical methanogens at Enceladus's  
250 seafloor consistent with methane emission flux in the space plume is  $1 \times 10^6 \text{ kg}_C \text{ yr}^{-1}$  [<sup>15,16</sup>].  
251 By multiplying this estimate by the biomass-specific demand of P for methanogens ( $\gamma_P$ ), a total  
252 P-demand by this hypothetical population of methanogens can be obtained. An upper estimate  
253 for  $\gamma_P$  is  $\approx 0.71 \text{ mmol}_P \text{ g}_{dw}^{-1}$  (Methods). Assuming that roughly half of a cell's dry weight  
254 is carbon<sup>17</sup>, the total demand of P by Earth-like methanogens at Enceladus's seafloor would  
255 be on the order of  $1 \times 10^5 \text{ mol}_P \text{ yr}^{-1}$ . Scaled to the mass of Enceladus's ocean (on the order  
256 of  $1 \times 10^{19} \text{ kg}$  [<sup>18</sup>]), this would result in a globally averaged consumption rate of the order of  
257  $10 \text{ nmol}_P \text{ L}^{-1} \text{ Myr}^{-1}$ . Geochemical calculations constrain bio-available P at concentrations  
258 exceeding  $100 \text{ nmolL}^{-1}$  [<sup>19,20</sup>]. Thus, methanogens could in principle exhaust the oceanic P

259 budget in as little time as 10 Myr. Geochemical calculations assume dissolution equilibrium,  
260 implying that depletion of dissolved P may allow more mineral P in the core to be dissolved.  
261 Most importantly, such a low concentration of P disagrees with analysis of plume composition<sup>21</sup>,  
262 which instead indicates concentrations of bio-available P in the ocean in excess of 1 mmolL<sup>-1</sup>,  
263 consistent with dissolution of a large fraction of P present in the core of chondritic composition<sup>19</sup>.  
264 Under this scenario, biotic consumption could be sustained for several billion years without  
265 noticeable impact on the concentration of P in the bulk ocean.

## 266 **1.9 Other bio-essential elements**

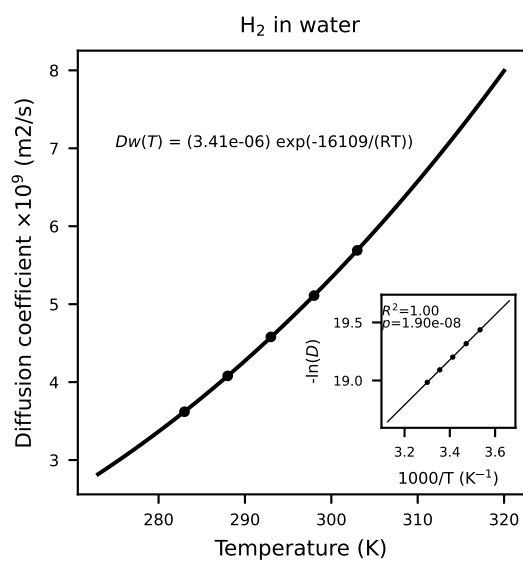
267 The availability of bio-essential nutrients other than P could constrain Enceladus's habitability  
268 or limit the growth of putative organisms. Nitrogen is expected to be abundant in Enceladus's  
269 ocean<sup>12</sup>, and recent findings support that Fe and S are also likely relatively abundant<sup>22</sup>. Deriving  
270 the growth thresholds for these elements requires that specific uptake parameters be estimated.  
271 However, knowledge of the order of magnitude for the biomass yields relative to these elements  
272 is sufficient to constrain whether biotic consumption may be commensurate with their estimated  
273 availability in Enceladus's ocean. As an approximation, it can be observed that at concentrations  
274 much greater than thresholds, the line in the 2-resource space that delimits trajectories that  
275 go to either one of two resource being limiting at biotic steady-state (equation 10 in Box 2 in  
276 Main text) becomes approximately one of line  $\gamma_2/\gamma_1$  with intersect 0.  
277 Interconversion of organic nitrogen and ammonia does not require the nitrogen atom to change  
278 oxidation state. Ammonia/ammonium can be used by microbes to assimilate nitrogen into  
279 biomass, and is returned by the degradation (including the fermentation) of N-containing  
280 organic matter, such as proteins. In anaerobic conditions, this would in principle constitute a  
281 simplified nitrogen cycle<sup>23</sup>. For the sake of sanity-checking the assumption that this cycling  
282 could keep methanogens in N-replete conditions, we roughly calculate the N demand implied  
283 by H<sub>2</sub>-limited primary production plausible for Enceladus's seafloor and compare this to the  
284 dissolved ammonium/ammonia stock implied by ammonia found in the plume<sup>12</sup>. Cleveland et  
285 al., 2007 [<sup>24</sup>] show the elemental ratio C:N:P (60:7:1) of soil microbial biomass to be globally  
286 constrained. The Redfield ratio suggests elemental ratios C:N:P closer to 106:16:1 in planktonic  
287 biomass. In the first order, this indicates that methanogens are expected to require between 1/7  
288 and 1/16 moles of N per mole of carbon that they fix. Replicating the same approach as in  
289 the previous, section, this indicates that a plausible nitrogen demand, scaled to ocean mass,  
290 for a population of H<sub>2</sub>-limited methanogens at Enceladus's seafloor be approximately 0.5

291 to  $1.2 \text{ nmolL}^{-1} \text{ kyr}^{-1}$  (see Table 1. in ref.<sup>25</sup> for the range of total ammonia concentration  
 292 assumed here). A credible range for ammonium/ammonia in Enceladus's ocean can be as low  
 293 as in the order of  $10 \text{ nmolL}^{-1}$  and up to  $\approx 1 \text{ mmolL}^{-1}$ , indicating that N-cycling in Enceladus's  
 294 seafloor would be required to turn nitrogen over on the 10,000 years timescale at the fastest,  
 295 longer than typical the typical timescale of biomass degradation. Thus, it does not appear an  
 296 implausible assumption that nitrogen cycling would sustain available nitrogen at non-limiting  
 297 concentrations. Second, at  $[\text{N}] \gg [\text{N}]^*$ , the line which separates  $\text{H}_2$ -limited steady-state from  
 298 N-limited steady states is  $[\text{N}] \approx (1/400) * [\text{H}_2]$  (taking  $\gamma_{\text{N}} \approx 1/10 \text{ molN per molC}$  and  $\gamma_{\text{H}_2} \approx 40$   
 299  $\text{molH}_2 \text{ per molC}$ ). This means that as an approximation, most of the plausible abiotic ranges  
 300 for  $\text{H}_2$  concentration and for N concentration satisfy  $[\text{N}] \geq (1/400) * [\text{H}_2]$ . In sum, N does  
 301 not seem to be a good candidate element for long term limitation of biomass production on  
 302 the long term in Enceladus's ocean, and while  $\text{H}_2$ -limited biotic steady-state appears more  
 303 plausible than N-limited biotic steady-state, the latter cannot be entirely ruled out.

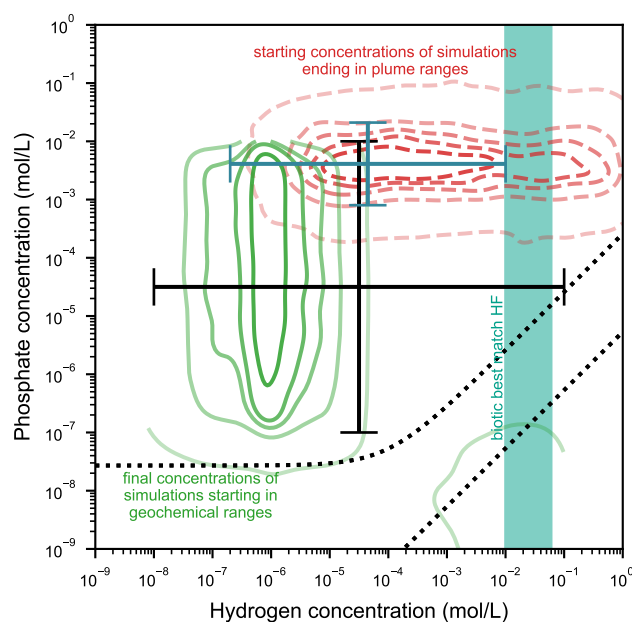
304 Xu et al., (2025) [<sup>22</sup>] show that iron (Fe) and sulfur (S) can be expected to be found at concentrations  
 305 in excess of  $[\text{Fe}] > 1 \times 10^{-8} \text{ molL}^{-1}$ ,  $[\text{S}] > 1 \times 10^{-6} \text{ molL}^{-1}$ . Both sulfur or iron, depending  
 306 on their oxidation state can be electron donor or acceptors of various metabolisms, but we  
 307 ignore this case here to focus on whether these elements could be sufficiently abundant to be  
 308 assimilated into the biomass of an ecosystem assumed to have methanogenesis as its primary  
 309 production performing metabolism. Xu et al. (2025) [<sup>22</sup>], compiled that the iron quota in  
 310 phytoplankton is in the order of  $10^{-4} \text{ mol per molC}$  or below, and the sulfur one is smaller  
 311 than  $3 \times 10^{-2} \text{ mol per molC}$ . Using the same approximation as above indicates that even at  
 312 the upper range for plausible abiotic  $\text{H}_2$  concentration from Glein and Truong (2025) [<sup>11</sup>],  
 313 around  $10 \text{ mmolL}^{-1}$ , and assuming the upper limit for cellular quotas, the concentration of Fe  
 314 needs to be greater than  $2.5 \times 10^{-8} \text{ molL}^{-1}$  and that of S greater than  $7.5 \times 10^{-6} \text{ molL}^{-1}$ . At  
 315 the highest,  $\text{H}_2$ -limited methanogens at Enceladus's seafloor would have a globally averaged  
 316 assimilation rate of S of  $0.25 \mu\text{molL}^{-1} \text{ Myr}^{-1}$ , and a assimilation rate of Fe of  $0.8 \text{ nmolL}^{-1} \text{ Myr}^{-1}$ .  
 317 Thus, limitation of growth by Fe or S is unlikely to occur for methanogens at Enceladus's  
 318 seafloor.

319 With sulfur, nitrogen, and phosphorus (and carbon and hydrogen), all major constituents known  
 320 biomass have been reviewed here. However, other minor elements could be proposed as limiting  
 321 factors for the growth of hypothetical methanogens in Enceladus's ocean. It is impossible, and  
 322 pointless, to consider all possible limiting elements, as the composition of hypothetical alien  
 323 life in Enceladus's ocean is unknown. However, it is interesting to briefly consider what minor  
 324 elements are typically found in the metabolic enzymes of Earth's methanogens to propose

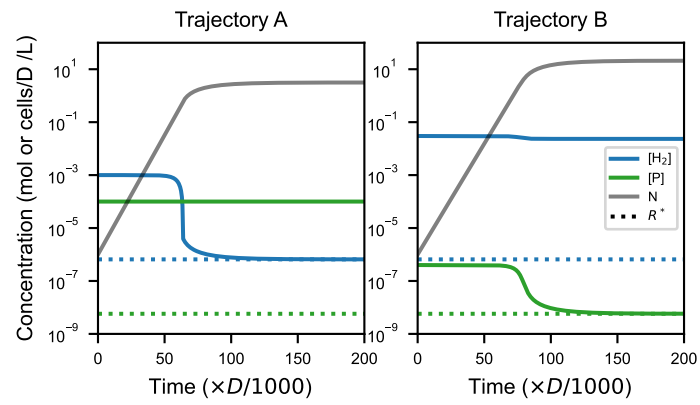
325 elements, the discovery of which in Enceladus's samples would comfort the hypothesis of  
326 the subsurface ocean being habitable to Earth-like methanogens. All known methanogens  
327 rely on ferredoxins, that is iron-sulfur proteins, for electron transfer steps of their metabolism.  
328 Because both iron and sulfur are expected in Enceladus's ocean, this should not pose an important  
329 additional constraint to our analysis above where the fraction of iron and sulfur in the bulk  
330 biomass is considered. Methanogenic archaea are also known to depend on the availability of  
331 nickel (Ni), and to a lesser extent on the availability of cobalt (Co) and molybdenum (Mo), for  
332 their growth<sup>26</sup>. Nickel is used in several of the key enzymes of methanogenic metabolisms,  
333 but may occasionally substituted for iron<sup>27</sup>. Ref.<sup>26</sup> have found that these metals no longer  
334 affect the growth of *M. thermoautotrophicum* when their concentration is greater than  $\approx 100 \text{ nmolL}^{-1}$ .  
335 Can these metals be expected to be available at concentrations  $> 100 \text{ nmolL}^{-1}$  in Enceladus's  
336 ocean? This depends primarily on the bulk composition of Enceladus's accretion materials  
337 and on whether these elements can be transferred to the ocean compartment at sufficient efficiency.  
338 Comets and chondrites have been suggested as having compositions analog to the accretion  
339 material of outer solar system bodies<sup>28</sup>, and silica nanograins found in Enceladus's plume  
340 have been suggested as evidence of hydrothermal alteration of a chondrite-like rocky core<sup>29</sup>.  
341 Ryugu samples were found to contain nickel, cobalt and molybdenum at appreciable levels<sup>30</sup>,  
342 and Bennu samples contain nickel and cobalt<sup>31</sup>. Moreover, Fisher and Ferrière (2026) [<sup>32</sup>]  
343 suggest the viability of the methanogenic metabolic network supported by elements present in  
344 asteroid and meteorites. Bulk abundance of cobalt, nickel and molybdenum in Ryugu samples  
345 suggest that if Enceladus's was accreted from similar materials, nickel and cobalt could be  
346 present in the ocean at concentrations greater than  $100 \text{ nmolL}^{-1}$ , and molybdenum at a concentration  
347 greater than  $10 \text{ nmolL}^{-1}$ , if entirely transferred to the ocean. The actual composition of Enceladus's  
348 accretionary material remains unknown, and the potential for these metals to leech into Enceladus's  
349 ocean is beyond the scope of this paper. Any hydrothermal setting at Enceladus's seafloor  
350 would likely be a hot spot for past or ongoing release of these metals into the ocean, hence  
351 we regard the scenario that methanogens be limited by the availability of these metals as less  
352 likely than the scenario of  $\text{H}_2$ -limitation.



**Figure S1 | Diffusion constant of H<sub>2</sub> in water as a function of temperature. Data from ref. <sup>8</sup>.**



**Figure S2 | Resource diagram with biotic initial and final states.** Error bars represent ranges of  $H_2$  and P concentrations inferred from chemical analysis of plume material (green) or derived from geochemical calculations (dark), they are the same as in Main Text Fig. 4. The green vertical span highlights the range of  $H_2$  concentrations for the hydrothermal fluid in the model of ref. <sup>14</sup> such that the  $H_2$  and  $CH_4$  fluxes escaping the hydrothermal system after consumption or production by hydrogenotrophic methanogens that best match escape rates in the plume (ref. <sup>15</sup>; *Supplementary text* section 1.7). The green levels correspond to the density of a sample of 50,000 simulations of the biotic steady-state resulting from initial states (log-)uniformly sampled in the concentration ranges derived from geochemical calculations (dark error bars). The red dashed levels show the density of starting conditions of simulations that have biotic steady-states that best match a 2D log-normal distribution centered on the intersection of the ranges for the inferred ocean composition from plume chemical analysis (green error bars).



**Figure S3** | Trajectories of Main text Fig. 4. as a function of time. Cell density has been rescaled to the dilution rate  $D$ , and the time by  $D/1000$  for legibility. Solid lines indicate the values of the concentrations of  $H_2$  (blue), P (green) and cells (gray). Dotted lines highlight the values of resource threshold  $[H_2]^*$  and  $[P]^*$  predicted under our 20-60 °C temperature theoretical scenario (Supplementary Table S1).

354 **Supplementary Tables**

**Table S1** | Summary of modeling parameterization scenarios.

	T range T point	0-20 C 0 C	20-60 C 40 C	60-120 C 80 C
Parameter	Unit			
$v_{max}^{H_2}$	mol/(hg <sub>dw</sub> )	0.1	1	10
$K_m^{H_2}$	molL <sup>-1</sup>	$5 \times 10^{-6}$	$2 \times 10^{-5}$	$6.6 \times 10^{-5}$
$\gamma_{H_2}$	mol <sub>H<sub>2</sub></sub> g <sub>dw</sub> <sup>-1</sup>	2.75	6.6	26.2
$v_{max}^P$	mol/(hg <sub>dw</sub> )	$1 \times 10^{-6}$	$1 \times 10^{-5}$	$1 \times 10^{-4}$
$K_m^P$	molL <sup>-1</sup>	$1 \times 10^{-8}$	$2.5 \times 10^{-8}$	$1 \times 10^{-7}$
$\gamma_P$	mol <sub>P</sub> g <sub>dw</sub> <sup>-1</sup>	$1.4 \times 10^{-4}$	$3.9 \times 10^{-4}$	$7.1 \times 10^{-4}$
$[H_2]^*$	molL <sup>-1</sup>	$2.75 \times 10^{-9}$	$6.54 \times 10^{-7}$	$5.6 \times 10^{-6}$
$[P]^*$	molL <sup>-1</sup>	$2.81 \times 10^{-11}$	$5.76 \times 10^{-9}$	$2.71 \times 10^{-8}$

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