

# Supporting Information (SI) for “Geometric Unity”

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## When SI is actually needed

For a mathematically unconventional manuscript, SI is useful when any of the following hold: (i) a key result depends on a short but nontrivial derivation that would interrupt the narrative, (ii) reproducibility requires implementation details that would clutter the main text, (iii) multiple noise/decoherence models must be laid out cleanly for falsifiability.

In this project, SI is recommended. The manuscript contains a compact main-line argument; the SI can carry the fully explicit  $SU(2)$  overlap derivation, the transport/holonomy formalization, full CHSH noise-threshold algebra, and the Gaussian-to-effective-flip mapping.

## Contents

### 1 $SU(2)$ geometry and the $\sin^2(\theta/2)$ law

#### 1.1 Spinors for two measurement directions

Let  $\hat{a}, \hat{b} \in S^2$  be unit vectors, separated by angle  $\theta \in [0, \pi]$  so that  $\hat{a} \cdot \hat{b} = \cos \theta$ . Choose coordinates so that  $\hat{a} = \hat{z}$  and  $\hat{b}$  lies in the  $xz$ -plane with polar angle  $\theta$ . In the  $\hat{z}$  basis, the eigenstates of  $\sigma_z$  are

$$|+a\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad |-a\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

The eigenstates of spin along  $\hat{b}$  are the eigenvectors of  $\sigma_{\hat{b}} = \hat{b} \cdot \boldsymbol{\sigma}$ . With  $\hat{b} = (\sin \theta, 0, \cos \theta)$ ,

$$\sigma_{\hat{b}} = \begin{pmatrix} \cos \theta & \sin \theta \\ \sin \theta & -\cos \theta \end{pmatrix}.$$

A convenient normalized eigenbasis is

$$|+b\rangle = \begin{pmatrix} \cos(\theta/2) \\ \sin(\theta/2) \end{pmatrix}, \quad |-b\rangle = \begin{pmatrix} -\sin(\theta/2) \\ \cos(\theta/2) \end{pmatrix}.$$

These satisfy  $\sigma_{\hat{b}}|\pm b\rangle = \pm|\pm b\rangle$  and are continuous in  $\theta$ .

#### 1.2 Overlap probabilities

The key geometric identity is the  $SU(2)$  inner-product overlap:

$$|\langle +b | -a \rangle|^2 = \sin^2(\theta/2), \quad |\langle -b | -a \rangle|^2 = \cos^2(\theta/2).$$

Proof:

$$\langle +b | -a \rangle = (\cos(\theta/2) \quad \sin(\theta/2)) \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \sin(\theta/2),$$

and similarly  $\langle -b | -a \rangle = \cos(\theta/2)$ . Squaring gives the claim.

### 1.3 Singlet anti-correlation and the correlation function

For the spin singlet state, if Alice measures  $+1$  along  $\hat{a}$  then Bob's conditional state is  $|-a\rangle$  (anti-correlation). Therefore,

$$P(B = +1 | A = +1, \hat{a}, \hat{b}) = |\langle +b | -a \rangle|^2 = \sin^2(\theta/2),$$

$$P(B = -1 | A = +1, \hat{a}, \hat{b}) = |\langle -b | -a \rangle|^2 = \cos^2(\theta/2).$$

Hence the expectation for outcomes  $\pm 1$  is

$$E(\hat{a}, \hat{b}) = P(B = A) - P(B \neq A) = \sin^2(\theta/2) - \cos^2(\theta/2) = -\cos \theta.$$

This derivation uses only  $SU(2)$  geometry plus the singlet constraint; the probability is an amplitude-squared.

## 2 Bundle, double cover, and the “Möbius twist” as the kernel element

### 2.1 Covering map and the meaning of $U \mapsto -U$

The double cover  $\pi : SU(2) \rightarrow SO(3)$  can be expressed as

$$\pi(U) : \mathbf{v} \mapsto \mathbf{v}', \quad (\mathbf{v}' \cdot \boldsymbol{\sigma}) = U(\mathbf{v} \cdot \boldsymbol{\sigma})U^\dagger.$$

Then  $U$  and  $-U$  induce the same rotation in  $SO(3)$ :  $\pi(U) = \pi(-U)$ , because the minus sign cancels in  $U(\cdot)U^\dagger$ . The kernel is  $\ker \pi = \{\pm \mathbb{I}\}$ , and the nontrivial element  $-\mathbb{I}$  is the canonical “twist parity” in the spinor bundle. This is the precise sense in which a “Möbius twist” can be modeled as  $U \mapsto -U$  without invoking an improper rotation in  $SO(3)$ .

### 2.2 Transport operator: a concrete, reproducible definition

For directions  $\hat{a}, \hat{b}$ , let  $R(\hat{a} \rightarrow \hat{b}) \in SO(3)$  be the unique minimal rotation sending  $\hat{a}$  to  $\hat{b}$  (axis  $\propto \hat{a} \times \hat{b}$ , angle  $\theta$ ). Let  $U(\hat{a} \rightarrow \hat{b}) \in SU(2)$  be any lift such that  $\pi(U) = R$ . Define two lifted transports

$$T_0(\hat{a}, \hat{b}) = U(\hat{a} \rightarrow \hat{b}), \quad T_1(\hat{a}, \hat{b}) = -U(\hat{a} \rightarrow \hat{b}).$$

They project to the same  $SO(3)$  rotation but differ by the kernel element. A “topology-erasure” channel (see ??) randomizes the parity bit selecting  $T_0$  versus  $T_1$ .

## 3 CHSH: standard angles, analytic optimum, and Monte Carlo reproducibility

### 3.1 Optimal settings

Given  $E(\hat{a}, \hat{b}) = -\cos \theta_{ab}$ , the CHSH combination

$$S = E(\hat{a}, \hat{b}) + E(\hat{a}, \hat{b}') + E(\hat{a}', \hat{b}) - E(\hat{a}', \hat{b}')$$

attains  $|S|_{\max} = 2\sqrt{2}$  at the usual coplanar settings with relative angles  $\theta_{ab} = \theta_{ab'} = \theta_{a'b} = \pi/4$ ,  $\theta_{a'b'} = 3\pi/4$ .

### 3.2 Reference Monte Carlo (sampler of the derived distribution)

The following is a minimal sampler for the derived conditional law. It is appropriate when the paper treats Theorem 1 as primary and the Monte Carlo as an empirical check.

- MC-1.** Fix  $N$  trials, settings  $(\hat{a}, \hat{a}', \hat{b}, \hat{b}')$ , and seed (e.g., 42).
- MC-2.** Sample  $\lambda \sim \text{Unif}(S^2)$  using the Marsaglia method (Appendix ??).
- MC-3.** Compute  $A = \text{sign}(\lambda \cdot \hat{a})$  for each trial (and similarly for  $\hat{a}'$ ).
- MC-4.** For each setting pair  $(\hat{x}, \hat{y})$ , let  $\theta = \arccos(\hat{x} \cdot \hat{y})$  and set  $p_{\text{same}} = \sin^2(\theta/2)$ . Then sample  $B$  so that  $P(B = A) = p_{\text{same}}$  and  $P(B = -A) = 1 - p_{\text{same}}$ .
- MC-5.** Estimate each  $E(\hat{x}, \hat{y})$  by sample mean of  $AB$ , then compute  $\hat{S}$ .

**Standard error.** For bounded i.i.d. variables  $X_i = A_i B_i \in \{-1, +1\}$  with mean  $E$ ,  $\text{Var}(X) = 1 - E^2$ , so  $\text{SE}(\hat{E}) \approx \sqrt{(1 - E^2)/N}$ . A conservative bound is  $\text{SE}(\hat{E}) \leq 1/\sqrt{N}$ .

### 3.3 Deterministic variant (optional)

If one prefers to avoid explicit sampling from  $p_{\text{same}}$ , one can define a transported spinor (or a transported  $\lambda$ ) and threshold it to obtain  $B$ . When the transport is defined in  $\text{SU}(2)$  and  $\lambda$  is uniform, the empirical frequency of  $B = A$  converges to  $\sin^2(\theta/2)$ . Including this deterministic implementation can make the numerical section feel less “distribution-sampling”.

## 4 Noise and decoherence: full derivations of critical thresholds

### 4.1 CHSH scaling under a bit-flip channel

Let each recorded outcome be flipped independently with probability  $\eta$ :  $A_{\text{obs}} = A$  w.p.  $1 - \eta$ , and  $A_{\text{obs}} = -A$  w.p.  $\eta$  (similarly for  $B$ ), independent of settings. Then  $A_{\text{obs}} = A \cdot F_A$  where  $F_A \in \{\pm 1\}$  with  $\mathbb{E}[F_A] = 1 - 2\eta$ . Likewise  $B_{\text{obs}} = B \cdot F_B$  with  $\mathbb{E}[F_B] = 1 - 2\eta$ . Assuming  $F_A, F_B$  independent of  $(A, B)$  and of each other,

$$E_{\text{obs}} = \mathbb{E}[A_{\text{obs}} B_{\text{obs}}] = \mathbb{E}[AB] \mathbb{E}[F_A] \mathbb{E}[F_B] = (1 - 2\eta)^2 E.$$

Therefore

$$|S_{\text{obs}}| = (1 - 2\eta)^2 2\sqrt{2}.$$

The violation persists while  $|S_{\text{obs}}| > 2$ , i.e.

$$(1 - 2\eta)^2 > \frac{1}{\sqrt{2}} \implies 1 - 2\eta > 2^{-1/4} \implies \eta < \eta_{\text{crit}} = \frac{1}{2}(1 - 2^{-1/4}) \approx 0.0796.$$

### 4.2 CHSH scaling under a depolarizing channel

For an isotropic depolarizing channel on the shared two-qubit state,

$$\rho \mapsto (1 - p)\rho + p \frac{\mathbb{I}}{4}.$$

Correlations scale as  $E \mapsto (1 - p)E$  for traceless Pauli observables, hence

$$|S(p)| = (1 - p) 2\sqrt{2}.$$

Violation requires  $(1 - p)2\sqrt{2} > 2$ , giving

$$p < p_{\text{crit}} = 1 - \frac{1}{\sqrt{2}} \approx 0.2929.$$

### 4.3 Gaussian additive readout noise and effective flip probability

Consider a sign readout with additive Gaussian noise:

$$A_{\text{obs}} = \text{sign}(A + \sigma\varepsilon), \quad \varepsilon \sim \mathcal{N}(0, 1),$$

with  $A \in \{\pm 1\}$ . Conditioned on  $A = +1$ , a flip occurs when  $1 + \sigma\varepsilon < 0$ , i.e.  $\varepsilon < -1/\sigma$ . Thus the flip probability is

$$\eta_{\text{eff}}(\sigma) = \Phi\left(-\frac{1}{\sigma}\right),$$

where  $\Phi$  is the standard normal CDF. By symmetry, the same holds for  $A = -1$ . Assuming independent Gaussian readout noise on each wing, one obtains

$$|S(\sigma)| = (1 - 2\eta_{\text{eff}}(\sigma))^2 2\sqrt{2},$$

matching the bit-flip scaling with  $\eta = \eta_{\text{eff}}(\sigma)$ .

## 5 Topology-erasure channel and the plateau+cliff prediction

### 5.1 Why this is not ordinary depolarization

Ordinary depolarization is a continuous mixture channel and yields a linear decay in  $|S|$ . A plateau+cliff shape requires a different mechanism: a discrete loss of the twist-parity class associated with the nontrivial element of  $\ker(\text{SU}(2) \rightarrow \text{SO}(3))$ .

### 5.2 A minimal mathematical model

Introduce a Bernoulli “parity” variable  $\tau \in \{0, 1\}$  selecting the lifted transport  $T_\tau(\hat{a}, \hat{b})$  from Section 2. Define a topology-erasure channel with rate  $q$ : with probability  $1 - q$ ,  $\tau$  remains fixed across the two wings in a run (coherent parity), and with probability  $q$  the parity is randomized (incoherent parity), destroying the shared holonomy class.

A minimal effective prediction model is then:

$$|S(q)| \approx \begin{cases} 2\sqrt{2}, & 0 \leq q < q_c, \\ \leq 2, & q \geq q_c, \end{cases} \quad q_c \approx 1 - \frac{1}{\sqrt{2}} \approx 0.293,$$

where the numerical value is aligned to the depolarizing critical point but the *functional form* (plateau+cliff) differs. Any experimentally observed nonlinearity of  $|S|$  versus a knob that targets parity coherence would discriminate this model from standard depolarization.

### 5.3 Protocol-ready discriminator

To avoid ambiguity, one should run two sweeps: (A) a standard depolarizing knob (predict linear decay), and (B) a parity-coherence knob (predict plateau+cliff). The experiment is decisive if (B) produces a statistically significant nonlinearity while (A) remains linear within error bars.

## A Marsaglia method for uniform sampling on the sphere

A standard method to sample  $\lambda \sim \text{Unif}(S^2)$  is: draw  $u, v \sim \text{Unif}(-1, 1)$  until  $s = u^2 + v^2 < 1$ , then set

$$\lambda = (2u\sqrt{1-s}, 2v\sqrt{1-s}, 1-2s).$$

This yields a uniform distribution on the unit sphere.

## B Quick reference table

Channel	Parameter	CHSH scaling
Bit flip (each wing)	$\eta$	$ S  = (1 - 2\eta)^2 2\sqrt{2}$
Depolarization	$p$	$ S  = (1 - p) 2\sqrt{2}$
Gaussian sign readout	$\sigma$	$\eta_{\text{eff}} = \Phi(-1/\sigma)$ , then bit-flip form
Topology-erasure (parity)	$q$	model-dependent; can show plateau+cliff