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2 **Supporting Information for**

3 **How urban heterogeneity and turbulence shape street-level heat exposure**

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7 **This PDF file includes:**

8 Supporting text

9 Figs. S1 to S18

10 Tables S1 to S3

11 SI References

12 **Supporting Information Text**

13 **A1. Derivation and Correspondence between  $\ell$  and  $\ell_\varepsilon$ .** In Santiago and Martilli (1), the dissipation mixing length  $\ell_\varepsilon$ , formulation  
 14 is defined via the relation  $C_k \ell = C_\mu \frac{\ell_\varepsilon}{C_\varepsilon}$ , where  $C_\mu = 0.09$ ,  $C_k = 0.4$ , and  $C_\varepsilon = 0.71$  (1). Their equations (16a–c) therefore  
 15 define  $\ell_\varepsilon$  (normalized by  $H$ ) as:

$$16 \quad \frac{\ell_\varepsilon}{C_\varepsilon H} = \begin{cases} \alpha'_1 \left(1 - \frac{d}{H}\right), & z \leq H \\ \alpha'_1 \left(\frac{z}{H} - \frac{d}{H}\right), & H < z \leq 1.5H \\ \alpha'_2 \left(\frac{z}{H} - \frac{d_2}{H}\right), & z > 1.5H \end{cases} \quad [A1]$$

17 To convert these expressions into a mixing length form  $\ell$ , one substitutes into the dimensional relationship above:

$$18 \quad \frac{\ell}{H} = C_\mu \frac{\ell_\varepsilon}{C_\varepsilon C_k} = \begin{cases} \alpha_1 \frac{C_\mu}{C_\varepsilon C_k} \left(1 - \frac{d}{H}\right), & z \leq H \\ \alpha_1 \frac{C_\mu}{C_\varepsilon C_k} \left(\frac{z}{H} - \frac{d}{H}\right), & H < z \leq 1.5H \\ \alpha_2 \frac{C_\mu}{C_\varepsilon C_k} \left(\frac{z}{H} - \frac{d_2}{H}\right), & z > 1.5H \end{cases} \quad [A2]$$

19 Thus, the coefficient  $\alpha_1$  in the classical mixing length formulation is given by:

$$20 \quad \alpha'_1 = \frac{C_\mu}{C_k} \alpha_1 = \frac{0.09}{0.71} \approx 0.225 \alpha_1 \quad [A3]$$

21 This yields a consistent mapping between the two formulations: reported values of  $\alpha'_1$  in the range 0.35–0.71 correspond  
 22 to  $\alpha_1 \approx 1.92$ –2.24, as used in the present study. Table S1 summarizes the reported values of  $\alpha_1$  and  $\alpha'_1$  along with the  
 23 corresponding values of  $\ell/H$  and  $\ell_\varepsilon/H$  as found in Santiago and Martilli (1), Simón-Moral et al. (2) and Nazarian et al. (3).

**Table S1. Reported values of mixing length parameters  $\alpha_1$  and  $\alpha'_1$  and the dissipation-based counterparts  $\ell/H$  and  $\ell_\varepsilon/H$ .**

Study	$\alpha_1$	$\alpha_2$	$\alpha'_1$	$\alpha'_2$
Santiago and Martilli (1) – RANS	2.24 – 2.25(2.24)	0.98 – 1.12(1.12)	0.504 – 0.506(0.504)	0.221 – 0.252(0.25)
Simón-Moral et al. (2014) – RANS	1.92 – 2.46(2.19)	1.14 – 1.27(1.2)	0.430 – 0.554(0.492)	0.257 – 0.286(0.27)
Nazarian et al. (2020) – LES	4.0	$\min(1.125, \max(0.45, 0.2925))$	0.5*	$\min(0.14, \max(0.056, 0.04))$ *

\*For consistency, a constant value of  $C_\mu$  was assumed. However, Nazarian et al. (2020) proposed a variable formulation for  $\alpha'_1$  as a function of  
 plan area density  $\lambda_p$  defined as  $C_\mu \begin{cases} \max(0.5, -\lambda^2 + 0.75\lambda_p + 0.022), & z \geq H \\ 0.05, & z < H \end{cases}$ .

24 **A2. Averaging Methodologies: Intrinsic vs. Comprehensive Averaging**

25 In building-resolving simulations of urban canopy flows, turbulent quantities must be horizontally averaged to derive representative  
 26 vertical profiles usable in urban canopy models. Two primary horizontal averaging approaches are:

- **Intrinsic averaging:** averages over the *fluid volume only*.
- **Comprehensive averaging:** averages fluid quantities but normalizes over the *total volume*, including buildings.

29 **A.2.1 Intrinsic Averaging.** The intrinsic horizontal average of a scalar variable  $\phi$  is defined as (4, 5):

$$30 \quad \langle \phi \rangle_I(z) = \frac{1}{A_a(z)} \int_{A_a(z)} \phi(x, y, z) dx dy \quad [A4]$$

31 where  $A_a(z)$  is the horizontal fluid area at height  $z$ .

32 Applying the spatial averaging theorem (5, 6), the derivative becomes:

$$33 \quad \left\langle \frac{\partial \phi}{\partial x_i} \right\rangle_I = \frac{\partial \langle \phi \rangle_I}{\partial x_i} + \langle \phi \rangle_I \frac{\partial \epsilon}{\partial x_i} + \frac{1}{A_a} \int_{\partial A_a} \phi n_i dl \quad [A5]$$

34 where  $\epsilon(z) = 1 - \lambda_p(z)$  is the local fluid fraction and  $n_i$  is the component of the unit normal vector at the fluid–solid interface.

35 **A.2.2 Comprehensive Averaging.** The comprehensive average is defined as:

36

$$\langle \phi \rangle_C(z) = \frac{1}{A} \int_{A_a(z)} \phi(x, y, z) dx dy \quad [A6]$$

37 with  $A = A_a(z) + A_p(z)$  the total horizontal area. The corresponding derivative is:

38

$$\left\langle \frac{\partial \phi}{\partial x_i} \right\rangle_C = \frac{\partial \langle \phi \rangle_C}{\partial x_i} + \frac{1}{A} \oint_{\partial A_a} \phi n_i dl \quad [A7]$$

39 Unlike intrinsic averaging, the comprehensive approach does not include a  $\partial \epsilon / \partial x_i$  term because the averaging volume does  
40 not vary with height.

41 **Implications and Application.**

42

- For **horizontal derivatives**, correction terms often vanish under horizontal homogeneity and no-slip conditions.

43

- For **vertical derivatives**, the correction terms are non-negligible, especially the  $\partial \epsilon / \partial z$  term in intrinsic averaging, which  
44 accounts for vertical changes in solid fraction.

45

- The interface integral accounts for discontinuities at building surfaces (e.g., pressure) and is generally zero for velocity  
46 under no-slip conditions.

47 In this study, we adopt **intrinsic averaging** for all velocity, flux, and turbulence analyses due to its clear physical  
48 interpretation and consistency with flow within the airspace, following (1, 4, 5).

49 **A2.1 Mixing-Length Formulations.** The turbulent momentum flux (TMF) is commonly modeled using a first-order closure:

50

$$\langle u' w' \rangle = -K_m \frac{\partial \langle u \rangle}{\partial z}, \quad \text{with} \quad K_m = \ell^2 \left| \frac{\partial \langle u \rangle}{\partial z} \right| \quad [A8]$$

51 Solving for  $\ell$ , we obtain the following formulations:

**Intrinsic Averaging.**

52

$$\ell_I = \sqrt{\left| \frac{\langle u' w' \rangle_I}{\frac{\partial \langle u \rangle_I}{\partial z} + \langle u \rangle_I \frac{\partial \epsilon}{\partial z}} \right|} \quad [A9]$$

**Comprehensive Averaging.**

53

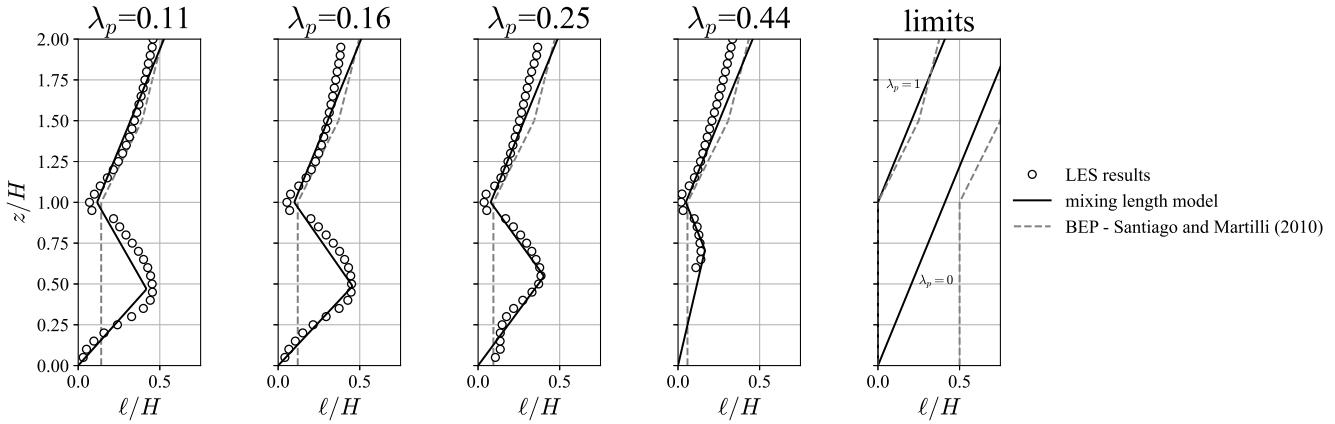
$$\ell_C = \sqrt{\left| \frac{\langle u' w' \rangle_C}{\epsilon \cdot \frac{\partial \langle u \rangle_C}{\partial z}} \right|} \quad [A10]$$

54 **A3. Turbulence mixing length closure**

55 **Mixing length prediction against LES results.** Figure S1 presents vertical profiles of the normalized mixing length ( $\ell/H$ ) derived  
56 from high-resolution large-eddy simulations (LES) for four urban canopies with varying plan area densities ( $\lambda_p = 0.11, 0.16,$   
57  $0.25, 0.44$ ). These LES results (open circles) serve as the reference dataset for evaluating the performance of the present model  
58 (solid lines) and the Building Effect Parameterization (BEP) by Santiago and Martilli (2010) (dashed gray lines). The rightmost  
59 panel shows the theoretical limiting cases at  $\lambda_p = 0$  and  $\lambda_p = 1$ , consistent with classical free-shear and fully obstructed-layer  
60 turbulence assumptions, respectively.

61 The present model demonstrates close agreement with the LES data across all densities, accurately capturing both the  
62 magnitude and vertical distribution of  $\ell/H$ . In contrast, the BEP formulation, while broadly consistent in trend, underpredicts  
63 mixing length magnitudes within the canopy layer and fails to represent the sharp transitions observed in the LES. The  
64 predictive fidelity of the new model is particularly evident near  $z/H = 1$ , where momentum exchange between canopy and  
65 overlying flow is most pronounced.

66 These results form the empirical foundation for the turbulence closure diagnostics and model development summarized in  
67 Figure 2 of the main text. By explicitly resolving the  $\ell/H$  structure across a range of  $\lambda_p$  using LES, we ensure the proposed  
68 formulation is rooted in high-fidelity data and applicable over the full range of realistic urban canopy conditions.



**Fig. S1.** Vertical profiles of normalized mixing length ( $\ell/H$ ) derived from large-eddy simulations (LES, circles), the present study's mixing length model (solid lines), and the Building Effect Parameterization (BEP) of Santiago and Martilli (2010) (dashed lines), for different plan area densities ( $\lambda_p$ ). The rightmost panel illustrates theoretical limit cases ( $\lambda_p = 0$  and  $\lambda_p = 1$ ).

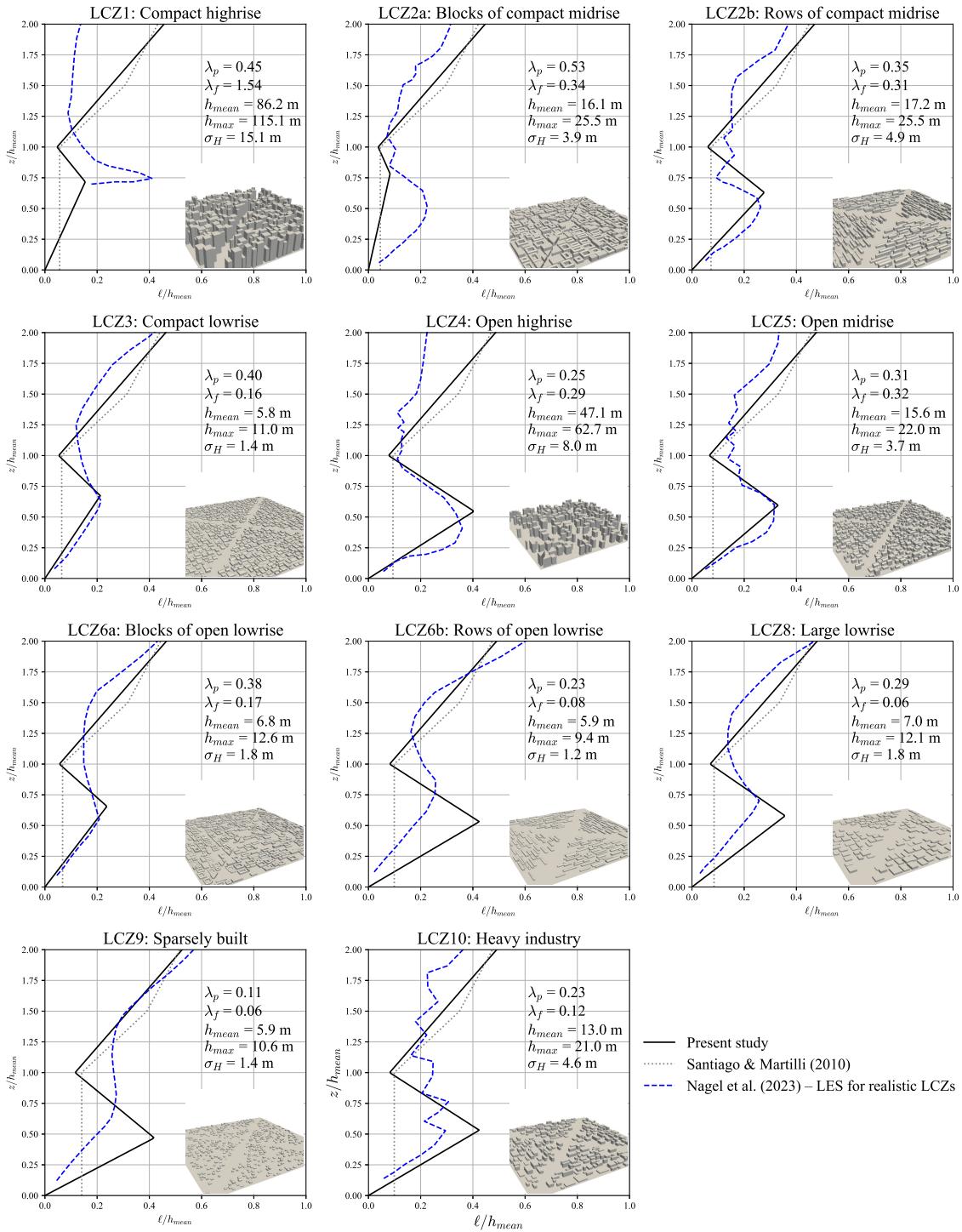
69 **Comparison of mixing length predictions against realistic and real world LES results.** To further evaluate the applicability of  
70 the proposed mixing-length formulation across realistic urban morphologies, we conducted comparisons against LES-derived  
71 profiles for both realistic Local Climate Zone (LCZ) classes (5) and real-world cities (7).

72 Figure S2 presents normalized vertical profiles of mixing length ( $\ell/h_{\text{mean}}$ ) for 12 LCZ categories simulated in Nagel et al.  
73 (2023) (5), capturing the breadth of global urban morphologies, from compact highrise (LCZ1) to sparsely built (LCZ9) and  
74 heavy industry (LCZ10). The present model (solid black lines) is benchmarked against the BEP formulation of Santiago and  
75 Martilli (2010) (dotted gray lines) and LES-derived profiles (dashed blue lines). Insets report the key morphological descriptors  
76 for each class, including plan area index ( $\lambda_p$ ), frontal area index ( $\lambda_f$ ), average and extreme building heights ( $h_{\text{mean}}$ ,  $h_{\text{min}}$ ,  
77  $h_{\text{max}}$ ), and height variability ( $\sigma_H$ ), with 3D representations of the building layouts. Across all typologies, the present model  
78 achieves stronger agreement with LES than the BEP scheme, accurately reproducing peak values and vertical structure of  $\ell$  in  
79 both dense and open configurations. This suggests that the closure is not limited to canonical canopies but is robust to diverse  
80 urban form factors.

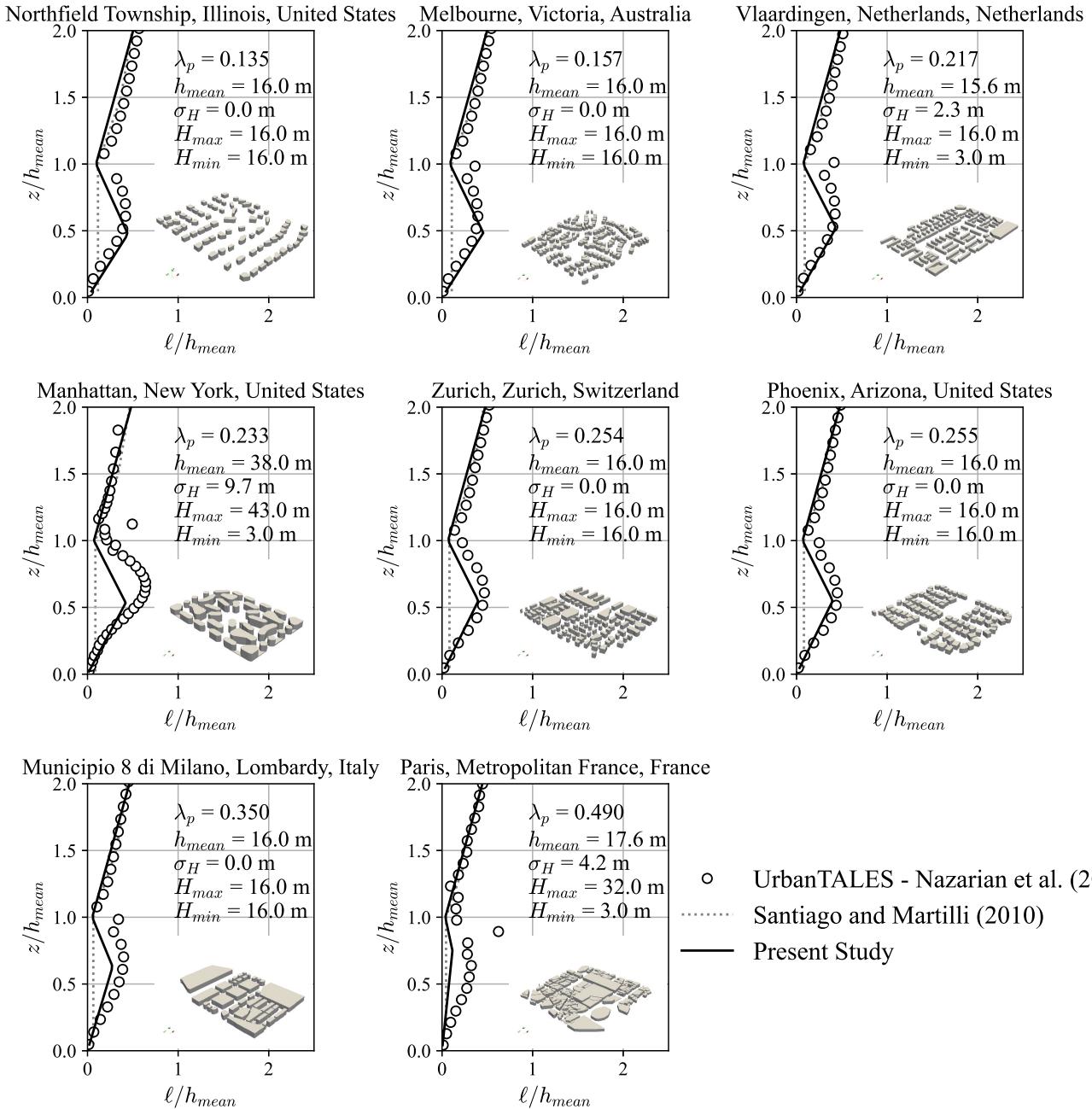
81 Figure S3 extends the analysis to nine real-world urban sites across North America, Europe, and Australia. Using LES  
82 results from the UrbanTALES database (Nazarian et al., 2025) (7), we compare the present model and BEP parameterization  
83 against  $\ell/h_{\text{mean}}$  profiles computed from fully resolved building-resolving LES. Despite variability in height distribution, layout  
84 regularity, and plan density across these cities, the present model consistently outperforms the Santiago and Martilli scheme. In  
85 particular, it reproduces the vertical extent and peak mixing values across complex layouts such as Manhattan and Paris, which  
86 include strong height gradients and nonuniform building arrangements. This supports the generalizability of the proposed  
87 formulation to operational applications involving realistic city geometries.

#### 88 **A4. Materials and methods: modeling tools**

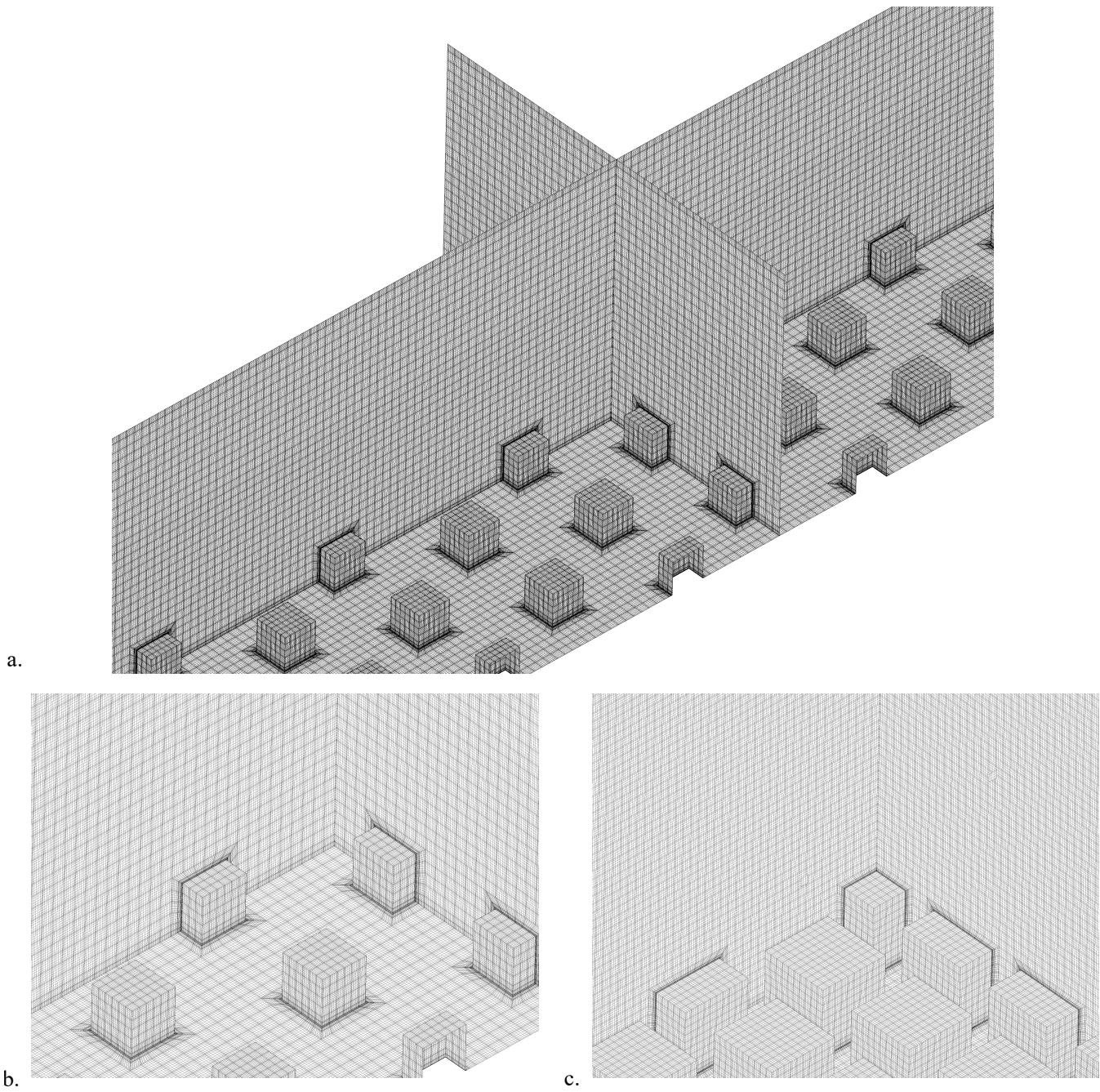
89 **Large Eddy Simulation using NekRS.** Using Large Eddy Simulation we expanded our data set via 3D simulations of canonical  
90 urban canopy flows using the high-order spectral element method (SEM) based computational fluid dynamics code **NekRS**  
91 (8). The computational domain was  $L_z = 8H$  tall while in the streamwise and spanwise dimensions  $L_x \times L_y$  the size of the  
92 domain was varied between  $L_x = 18H$  and  $30H$  for  $\lambda_p = 0.11$  and  $0.44$  respectively, The spanwise direction was  $L_y = 12$   
93 for most of the cases except  $\lambda_p = 0.16$  where the spanwise dimension of the domain was  $L_y = 12.5$ . The resolution of the  
94 SEM elements grid was  $0.25H$  for the cases with  $\lambda_p \leq 0.25$  ( $\lambda_p = 0.11, 0.16, 0.25$ ), which combined with the order 7th which  
95 corresponds to 8 Gauss-Lobatto-Legendre points gives a resolution of  $(0.25/8)H = 0.03125H$  (32 points per  $H$ ). For the case  
96 of  $\lambda_p = 0.44$ , the element size was reduced to  $0.125H$  and the polynomial order used was 5th order which corresponds to 6  
97 Gauss-Lobatto-Legendre points gives a resolution of  $(0.125/6)H = 0.0208H$  (48 points per  $H$ ). It is important to stress that  
98 an additional E element was added near the ground and cuboid surface (see fig. S4). A full summary of all simulated cases and  
99 their associated morphological parameters is provided in Table S2.



**Fig. S2.** Vertical profiles of normalized mixing length ( $\ell/h_{\text{mean}}$ ) for 12 realistic Local Climate Zone (LCZ) classes based on high-fidelity large-eddy simulations (LES) Nagel et al. (2023) (5). Each panel compares results from the present study's mixing length model (solid black lines) with the Building Effect Parameterization (BEP) of Santiago and Martilli (2010) (dotted gray lines) and LES-based profiles reported by Nagel et al. (2023) (dashed blue lines). Insets provide key morphological parameters: plan area index ( $\lambda_p$ ), frontal area index ( $\lambda_f$ ), mean height ( $h_{\text{mean}}$ ), the minimum and maximum building heights ( $h_{\text{min}}$ ,  $h_{\text{max}}$ ) and the standard deviation of heights ( $\sigma_H$ ). A 3D rendering of the corresponding building layout is shown in the lower-right corner of each panel. The profiles highlight the variability in mixing behavior across different urban typologies and demonstrate the improved agreement of the present model with LES data over conventional parameterizations.



**Fig. S3.** Vertical profiles of normalized mixing length ( $\ell/h_{mean}$ ) for selected real-world urban morphologies across the United States, Europe, and Australia. Each panel shows comparisons between LES results from UrbanTALES by Nazarian et al. (2025) (open circles) (7), the present study's mixing length model (solid black lines), and the Building Effect Parameterization (BEP) of Santiago and Martilli (2010) (dotted gray lines). Urban domains span a range of plan area densities ( $\lambda_p$ ), mean and maximum building heights ( $h_{mean}$ ,  $H_{max}$ ), and the heights standard ( $\sigma_H$ ). Each panel includes a 3D rendering in the lower-right corner, illustrating the associated building arrangement. The results demonstrate that the present model better captures both the shape and magnitude of the mixing length profiles across diverse urban layouts, outperforming conventional parameterizations especially in cases with height variability and non-uniform building spacing.



**Fig. S4.** Computational grid used for the numerical simulations of urban canopy flows, showing (a) an overview of the grid arrangement ( $\lambda_p = 0.11$ ), (b) a detailed view near building corners ( $\lambda_p = 0.11$ ), and (c) a close-up view of the refined grid resolution near building surfaces ( $\lambda_p = 0.44$ ).

**Table S2. Geometric and flow parameters for present LES and reference urban canopy simulations used for comparative analysis and mixing length model development.**

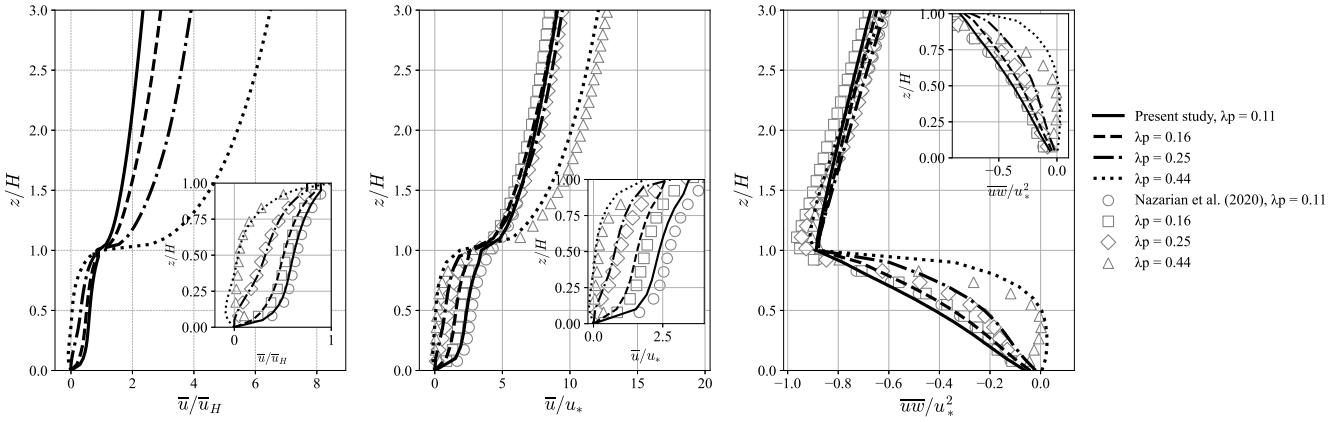
Present study – LES									
Dataset	$\lambda_p - S/A$	$(L_x \times L_y \times L_z)/H$	$N_x$	$N_y$	npoints	$\text{Re}_\infty = \frac{u_\infty H}{\nu}$	$\text{Re}_H = \frac{u_H H}{\nu}$	$\text{Re}_\tau = \frac{u_* H}{\nu}$	$\text{Re}_b = \frac{u_b H}{\nu}$
Present study	0.11 – $S$	$30 \times 12 \times 8$	10	4	98,058,240	7,589	2,683	723.4	6,000
	0.16 – $S$	$25 \times 12.5 \times 8$	10	5	87,654,400	7,336	2,286	736.0	6,000
	0.25 – $S$	$20 \times 12 \times 8$	10	6	70,901,760	6,854	1,891	616.0	6,000
	0.44 – $S$	$18 \times 12 \times 8$	12	8	204,456,960	6,539	1,534	487.8	6,000
Literature studies									
Dataset	$\lambda_p - S/A$	$(L_x \times L_y \times L_z)/H$	$N_x$	$N_y$		$\text{Re}_\infty = \frac{u_\infty H}{\nu}$			
Nazarian et al. (2020) (3)	0.065 – $S$	$24 \times 12 \times 7.4$	6	3		$10^6$			
	0.11 – $S$	$18 \times 9 \times 7.4$	6	3		$10^6$			
	0.16 – $S$	$15 \times 7.5 \times 7.4$	6	3		$10^6$			
	0.25 – $S$	$12 \times 6 \times 7.4$	6	3		$10^6$			
	0.44 – $S$	$10 \times 5 \times 7.4$	6	3		$10^6$			
	0.065 – $A$	$20 \times 12 \times 7.4$	5	3		$10^6$			
	0.44 – $A$	$7.5 \times 4.5 \times 7.4$	5	3		$10^6$			
	0.33 – $A$	$12 \times 9 \times 8$	6	3		6,500			
Castro et al. (2017) (9)	0.33 – $A'$	$12 \times 12 \times 12$	6	4		7,500			
	0.25 – $A$	$16 \times 16 \times 8$	8	8		4,750			
Branford et al. (2011) (10)	0.25 – $A'$	$16 \times 16 \times 8$	8	8		4,750			
	0.25 – $S$	$16 \times 16 \times 10$	8	8		4,790			
Xie et al. (2008) (11)									

\*Xie et al. (2008) (11) performed simulations at  $\lambda_p = 0.25$  using staggered cuboids with equal length and width but varying heights. Within each 4-by-4 unit, the building heights consisted of  $1 \times 0.26H$ ,  $3 \times 0.63H$ ,  $7 \times 0.98H$ ,  $4 \times 1.33H$ , and  $1 \times 1.68H$ , where  $H$  is the average building height.

100 The  $P_N - P_N$  spectral element formulation was adopted for all simulations (12). In this formulation, the velocity and pressure  
101 fields are both represented using tensor-product polynomials of degree  $p$  within each non-overlapping hexahedral element  $\Omega_e$ ,  
102 such that the global domain is  $\Omega = \bigcup_{e=1}^E \Omega_e$ . The basis functions are Lagrange interpolants defined on Gauss–Lobatto–Legendre  
103 (GLL) quadrature points, which ensures numerical stability and enables efficient pointwise quadrature operations within each  
104 element.

105 The incompressible filtered Navier–Stokes equations were solved using the GPU-accelerated spectral element solver NekRS,  
106 which provides exponential spatial convergence and second-order temporal accuracy. A de-aliasing procedure was applied to the  
107 nonlinear advection terms using the standard 3/2 over-integration rule (12). To balance spatial accuracy and computational  
108 performance, seventh-order elements  $p = 7$  ( $N = 8$  number of points per element) were used throughout. Time integration  
109 followed a mixed implicit–explicit scheme: all linear terms (e.g., viscous diffusion, pressure gradient) were treated implicitly  
110 using a projection method with pressure–velocity decoupling, while the nonlinear advection terms were treated explicitly. The  
111 time advancement used second-order backward differentiation (BDF2) for implicit terms and third-order extrapolation (EXT3)  
112 for explicit terms, consistent with a BDF2/EXT3 scheme. The LES was based on a spectral filtering approach, which act  
113 as an implicit high-pass filter for unresolved turbulent scales (13–15), thereby avoiding the need for an explicit subgrid-scale  
114 modeling using a transport equation.

115 Figure S5 presents a validation of the present study’s large-eddy simulation (LES) results through a comparison of vertical  
116 profiles of mean velocity and Reynolds shear stress with previously published data. The panels show normalized mean velocity  
117 profiles ( $\bar{u}/\bar{u}_H$  and  $\bar{u}/u_*$ ) and Reynolds shear stress ( $\bar{u}'\bar{w}'/u_*^2$ ) across four different urban canopy densities ( $\lambda_p = 0.11$ ,  $0.16$ ,  
118  $0.25$ , and  $0.44$ ). Simulation results from the present study are shown as lines, while symbols represent benchmark LES data  
119 from Nazarian et al. (2020) (7), enabling direct comparison across cases. The agreement between the two independent datasets  
120 confirms the fidelity of the current simulations in capturing both the velocity structure and momentum fluxes in and above  
121 the canopy. In particular, the present results match both the shape and magnitude of normalized velocity and stress profiles,  
122 reproducing key trends such as the shift of the velocity maximum with increasing  $\lambda_p$  and the inflection near the canopy top.  
123 Insets highlight the region inside the canopies.



**Fig. S5.** Vertical profiles of normalized mean velocity ( $\bar{u}/\bar{u}_H$  and  $\bar{u}/u_*$ ) and Reynolds shear stress ( $\bar{u}\bar{w}/u_*^2$ ) for different plan area densities ( $\lambda_p$ ), comparing the present study (lines) with numerical results from Nazarian et al. (2020) (symbols). Insets provide detailed views within the urban canopy.

Figure S6 presents one-dimensional turbulence spectra at multiple heights above urban canopies for a range of plan area densities ( $\lambda_p$ ). Left panels show normalized velocity spectra  $E(k_x H)/u_{*H}^2$  plotted against nondimensional wavenumber  $k_x H$ , with the characteristic  $-5/3$  slope of the inertial subrange indicated with an orange line. Right panels show premultiplied spectra  $k_x H E(k_x H)/u_b^2 H$  as a function of normalized wavelength  $\lambda_x/H$  at selected vertical locations ( $z/H = 1.05, 1.2, 1.5, 2.0$ ). The red dashed line marks the effective resolution limit, defined as twice the element size ( $2\Delta$ ).

The spectra confirm that all cases are well resolved in the inertial subrange. A clear  $-5/3$  slope is observed across approximately one decade in  $k_x H$ , and the absence of energy accumulation near the resolution cutoff indicates low numerical dissipation and accurate capture of turbulent scales. Moreover, the energy content associated with scales smaller than the filter size is several orders of magnitude lower than that in the inertial range. The premultiplied spectra reveal changes in the distribution of turbulent energy with height. Near the canopy top ( $z/H \approx 1.05$ ), energy is distributed more broadly across wavelengths due to strong wake-induced turbulence and multi-scale interactions. At higher elevations, the spectra become more sharply peaked. In all cases, the energy carried by scales near the filter size remains negligible, further confirming that the simulations are well resolved.

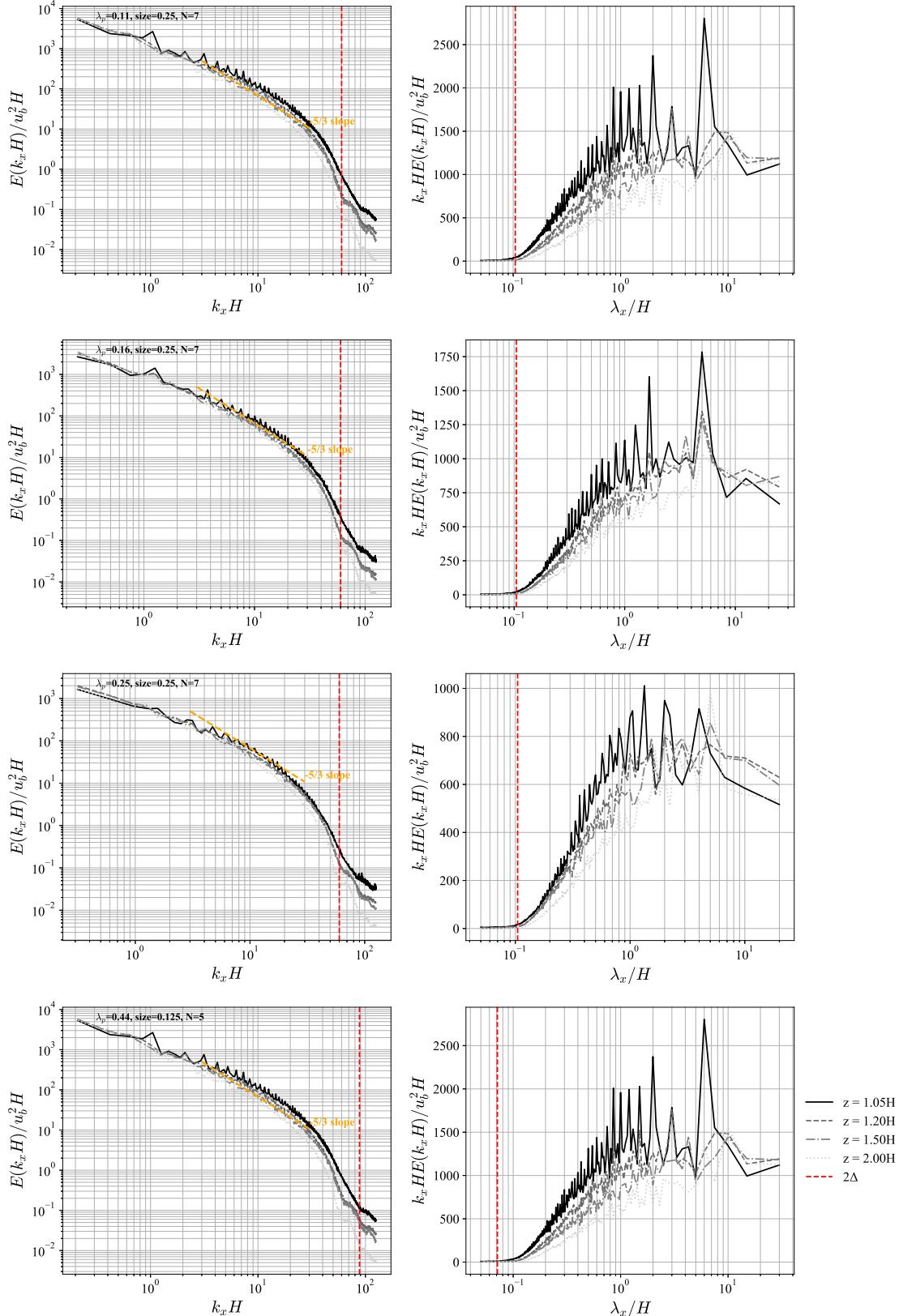
**WRF Model Setup and Experiment Design.** In this study, we used version 4.5.1 of the Weather Research and Forecasting (WRF) model (16), a non-hydrostatic, fully compressible numerical weather prediction system. To represent urban processes, we incorporated the Building Effect Parameterization (BEP) and the Building Energy Model (BEM) (17–19), which simulate urban canopy layer dynamics and associated building energy exchanges.

The simulations employed three two-way nested domains with horizontal grid spacings of 4.5 km, 1.5 km, and 500 m, respectively. The outer domain consisted of  $404 \times 403$  grid points, the intermediate domain of  $412 \times 412$  grid points, and the innermost domain—covering the Chicago Metropolitan Area (CMA)—also used  $412 \times 412$  grid points (Figure S7a). The model integration time step was set to 15 seconds and was further adjusted for each domain based on Courant–Friedrichs–Lowy (CFL) conditions.

The vertical grid included 55 eta levels extending from the surface to approximately 20 km in altitude. To adequately capture lower atmospheric dynamics, enhanced vertical resolution was used near the surface, with 13 levels below 100 m and the first 5 layers located within 10 m above ground, each having a resolution of approximately 1–2 m. BEP-based models employ a secondary “urban grid” to resolve sub-grid urban morphology. For this study, we set the maximum number of vertical levels in the urban grid to 30, the vertical resolution of the urban grid to 3.0 m, and the maximum number of building types per urban class to 20.

Initial and boundary conditions were obtained from the fifth-generation ECMWF Reanalysis (ERA5) dataset (20), providing 3-hour updates at a spatial resolution of approximately 31 km. Simulations were initialized using ERA5 atmospheric fields, with land surface and soil variables interpolated to the WRF grid. To improve model accuracy over the Great Lakes region, ERA5 sea surface temperature (SST) fields were replaced with higher-resolution observations from the NOAA Great Lakes Surface Environmental Analysis (GLSEA) product (21), which provides daily maps at 1.3 km resolution. This substitution has been shown to enhance WRF performance in simulating near-surface air temperature and relative humidity over the Great Lakes (22).

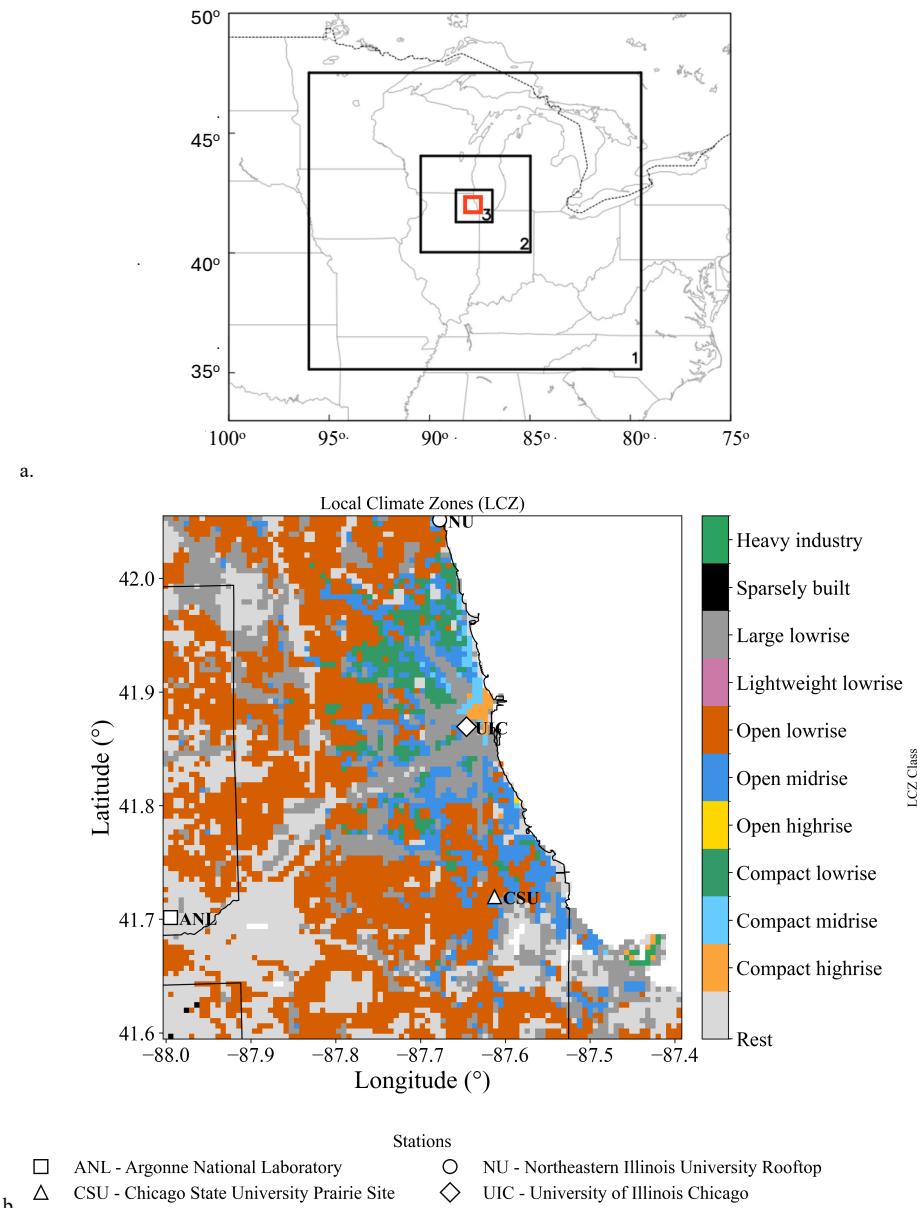
The physics suite included the Noah land surface model (23), the Bougeault–Lacarrère planetary boundary layer (PBL) scheme (24), WRF Single-Moment 6-class microphysics (25), the Dudhia shortwave radiation scheme (26), and the Rapid Radiative Transfer Model (RRTM) for longwave radiation (27). WRF Preprocessing System (WPS) tools 4.6.0 were used for initial and boundary condition preparation. The simulations were initialized at 18 UTC on July 21, 2024, and ran until 00 UTC on July 24, 2024. The first 6 hours were discarded to allow for model spin-up. Model outputs were archived at hourly intervals. Urban geometry was characterized using Local Climate Zones (LCZs) derived from the World Urban Database and Access Portal Tool (WUDAPT) (28), and enhanced with high-resolution data from the Height Mapping of Trees and Buildings



**Fig. S6.** One-dimensional turbulence spectra at various heights above urban canopies with different plan area densities ( $\lambda_p$ ). Left panels show normalized spectra  $E(k_x H) / u_b^2 H$  plotted against wavenumber  $k_x H$ , with the characteristic  $-5/3$  inertial subrange slope indicated. Right panels display premultiplied spectra  $k_x H E(k_x H) / u_b^2 H$  versus normalized wavelength  $\lambda_x / H$  at different vertical locations ( $z/H$ ). The red dashed vertical line indicates twice the grid spacing ( $2\Delta$ ), marking the limit of spatial resolution.

166 (HiTAB) dataset (29). Within the Chicago Metropolitan Area (CMA), the urban geometry was further refined using the  
 167 University of Texas – Global Building Heights for Urban Studies (UT-GLOBUS) dataset (30). Eleven distinct LCZ classes  
 168 were defined based on prior studies and spatial datasets (31). Albedo and emissivity distributions over the Chicago area were  
 169 refined using high-resolution data from Cheng et al. (32). Figure S7b shows the spatial distribution of LCZs applied over  
 170 domain 3 (at a 500m resolution), with the red box in Figure S7a indicating the extent shown in Figure S7b.

171 To evaluate the robustness and added value of the proposed turbulence model, we conducted four comparative simulations  
 172 over the CMA, varying both the urban morphology dataset and the turbulence parameterization. This design allows us  
 173 to disentangle the respective influences of improved turbulence representation and enhanced urban morphology on model  
 174 performance. The two levels of morphological detail—WUDAPT-based LCZs and the high-resolution HiTAB dataset—were  
 175 each paired with both the original Santiago and Martilli (2010) scheme (1) and the new model proposed in this study. This  
 176 setup enables a side-by-side evaluation of model skill across combinations of physical parameterizations and morphology  
 177 resolutions, highlighting their respective roles in capturing spatial heterogeneity in urban microclimate conditions.



**Fig. S7.** (a) WRF domain. The red box indicates the Chicago Metropolitan Area (CMA). (b) The urban land-use categories after incorporating local climate zones (LCZs) from WUDAPT data over domain 3.

178 **A5. WRF data analysis**

179 **Model validation.** To evaluate model performance against in-situ observations, we selected four urban monitoring sites within  
 180 the Chicago Metropolitan Area, shown in Fig. S7b (33, 34). These sites span a range of urban contexts and LCZs, and  
 181 include rooftop and near-surface installations. Table S3 summarizes their geographic coordinates and measurement heights  
 182 above ground level. Each of these sites is equipped with a Vaisala Weather Transmitter (WXT530 series), which provides  
 183 measurements of ambient temperature ( $\pm 0.3^\circ\text{C}$  accuracy), pressure ( $\pm 0.5 \text{ hPa}$  at  $0 \dots +30^\circ\text{C}$  accuracy), and relative humidity  
 184 ( $\pm 3\%$  RH at RH  $< 90\%$  accuracy) via an RC oscillator and dual reference capacitors within the sensor PTU module.

185 Figures S8-S10 present time series comparisons of modeled and observed meteorological variables at the four urban monitoring  
 186 sites at Argonne National Laboratory (ANL), Chicago State University (CSU), Northwestern University (NU), and University  
 187 of Illinois Chicago (UIC) within the CMA. Each panel shows results for the present study and the Santiago and Martilli (2010)  
 188 (1) scheme, using both WUDAPT and HiTAB urban morphology datasets, against observational data (open circles).

Table S3. Station metadata used for model–observation comparison.

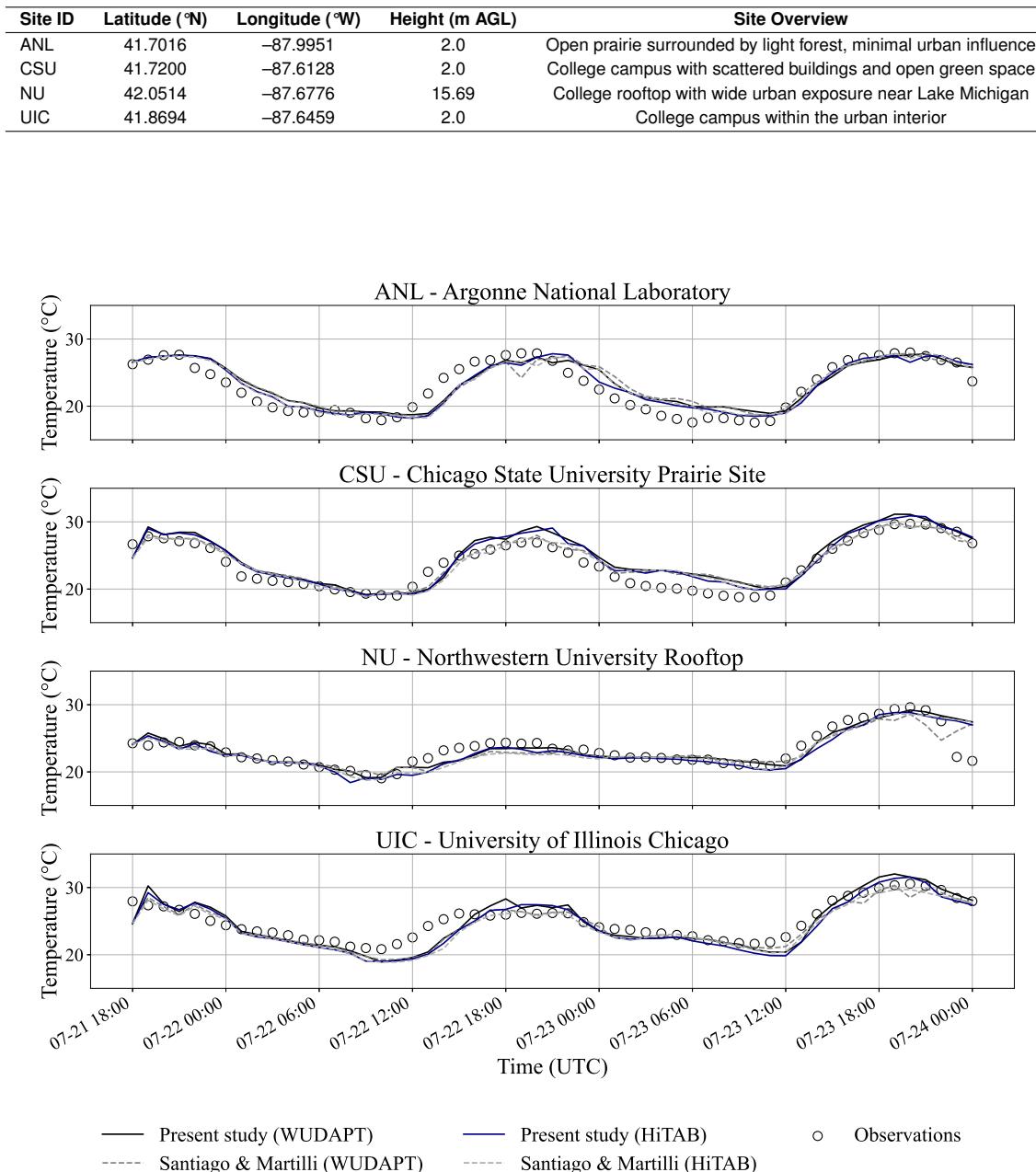


Fig. S8. Time series of observed and modeled temperature at four urban observation sites (ANL, CSU, NU, and UIC) from July 21 to July 24 2024. The present study's simulations using WUDAPT and HiTAB urban morphology datasets are shown with black and blue lines, respectively, and are compared to the Santiago and Martilli (2010) scheme (red and gray dashed lines). Observations are shown as open circles.

189 Figure S8 shows near-surface air temperature, revealing consistent diurnal patterns and generally good agreement between  
190 modeled and observed temperatures. The present study tends to better capture daytime peaks and nighttime minima,  
191 particularly at NU and UIC. Figure S9 displays near-surface relative humidity (RH), where the present model generally  
192 improves both the magnitude and timing of RH evolution, especially during day–night transitions. This improvement is most  
193 evident at ANL and CSU, with good agreement also observed at NU and UIC despite some overestimation during nighttime  
194 hours. Figure S10 shows wind speed comparisons, with larger variability and site-specific differences, though the present study  
195 captures temporal trends effectively. Further quantitative evaluation of these time series, including regression, RMSE, % error  
196 and observed-versus-modeled scatter plots, is provided in the main text (see Fig. 4), where we further assess the predictive  
197 performance of each configuration.

198 **Relative Humidity and Cooling Energy Consumption**. Figures S11 and S12 provide four experiment simulated statistical  
199 comparisons of RH and urban cooling energy consumption between the Present Study and the Santiago and Martilli (2010)  
200 scheme (S&M) across the CMA.

201 Figure S11 shows RH distributions from both models and datasets (panels a–b, e–f), along with spatial maps of  $\Delta$ RH (Present  
202 – S&M; panels c, g) and associated probability density functions (panels d, h). The present model yields lower RH values  
203 in regions that exhibit higher temperatures (see Fig. 4), consistent with increased turbulent mixing. While domain-averaged  
204 RH differences are modest (typically 1–6% for WUDAPT and 1–3% for HiTAB), the spatial coherence of  $\Delta$  RH patterns  
205 reflects canopy-layer sensitivity to updated turbulence representations. Inter-dataset differences (HiTAB – WUDAPT; panels  
206 i–j) further reveal the role of urban morphology in modulating humidity. Violin plots (panels k–l) summarize the effect of  
207 turbulence modeling approaches on the full RH distributions and  $\Delta$  RH across all cases. The present study generally shows  
208 higher RH values and distinct distribution shapes compared to S&M, with differences in spread and central tendency reflecting  
209 sensitivity to both canopy formulation and urban morphology dataset.

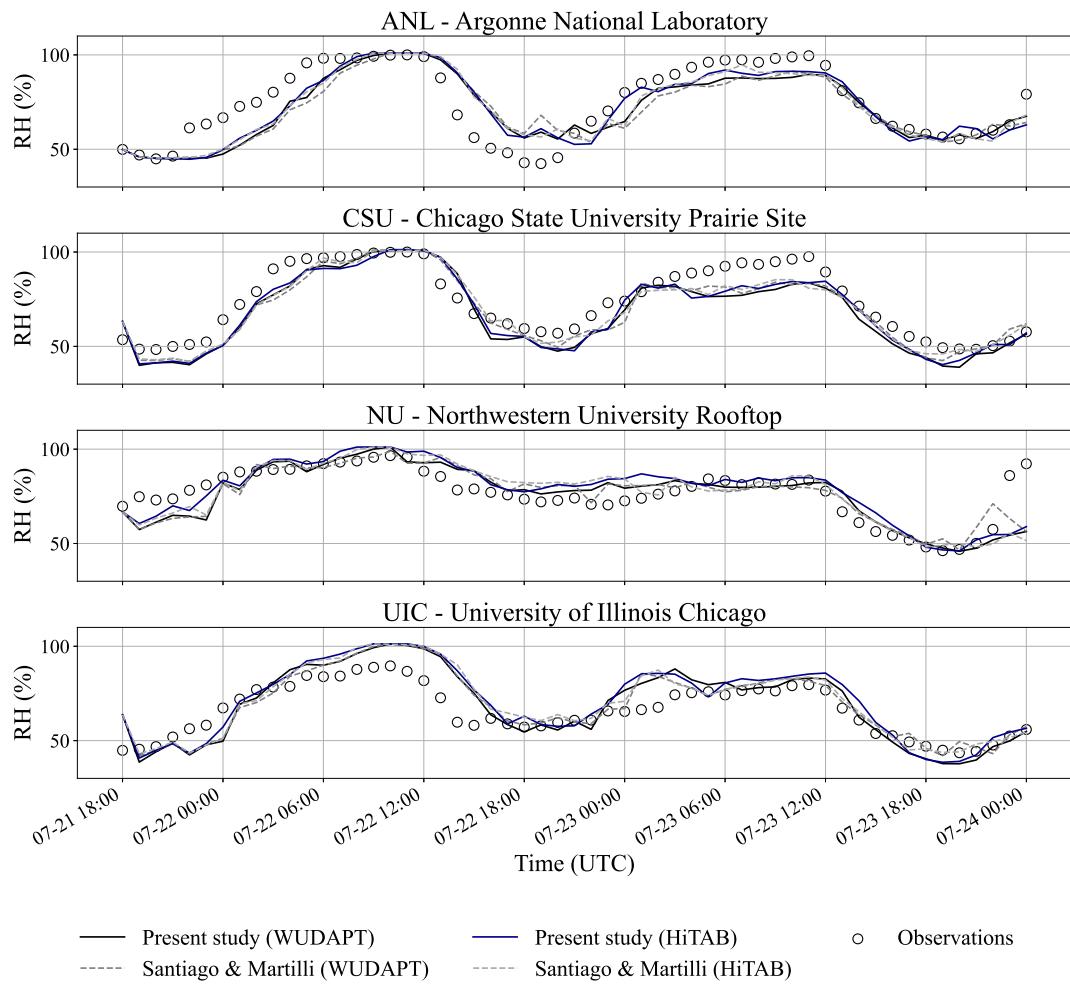
210 Figure S12 provides a similar comparison for building energy consumption. Panels a–b and e–f show total cooling energy  
211 use from the Present Study and S&M, while panels c and g display the spatial differences ( $\Delta$  Energy = Present – S&M)  
212 along with corresponding probability density functions. Energy use is highly sensitive to both the urban canopy representation  
213 and the turbulence parameterization: the present model systematically predicts higher cooling demand, especially in denser  
214 areas where canopy-layer mixing is more strongly enhanced. Although the absolute values are larger in WUDAPT due to its  
215 greater total building volume (resulting from a higher number of floors), the overall patterns and inter-model differences remain  
216 consistent across both morphology datasets. Inter-dataset contrasts (panels i–j) and violin plots (panels k–l) further reinforce  
217 these findings, with the Present Study showing higher mean, median, and maximum energy values, and broader distribution  
218 tails relative to S&M.

219 **LCZ-Resolved temperature, RH, and Cooling energy distributions.** Figures S13–S18 present LCZ-resolved probability distributions  
220 of air temperature, RH, and cooling energy consumption across Cook County, Illinois, using both WUDAPT and  
221 HiTAB urban morphology datasets. These distributions provide a detailed evaluation of model behavior within distinct urban  
222 typologies and complement the domain-wide analyses discussed in the main text.

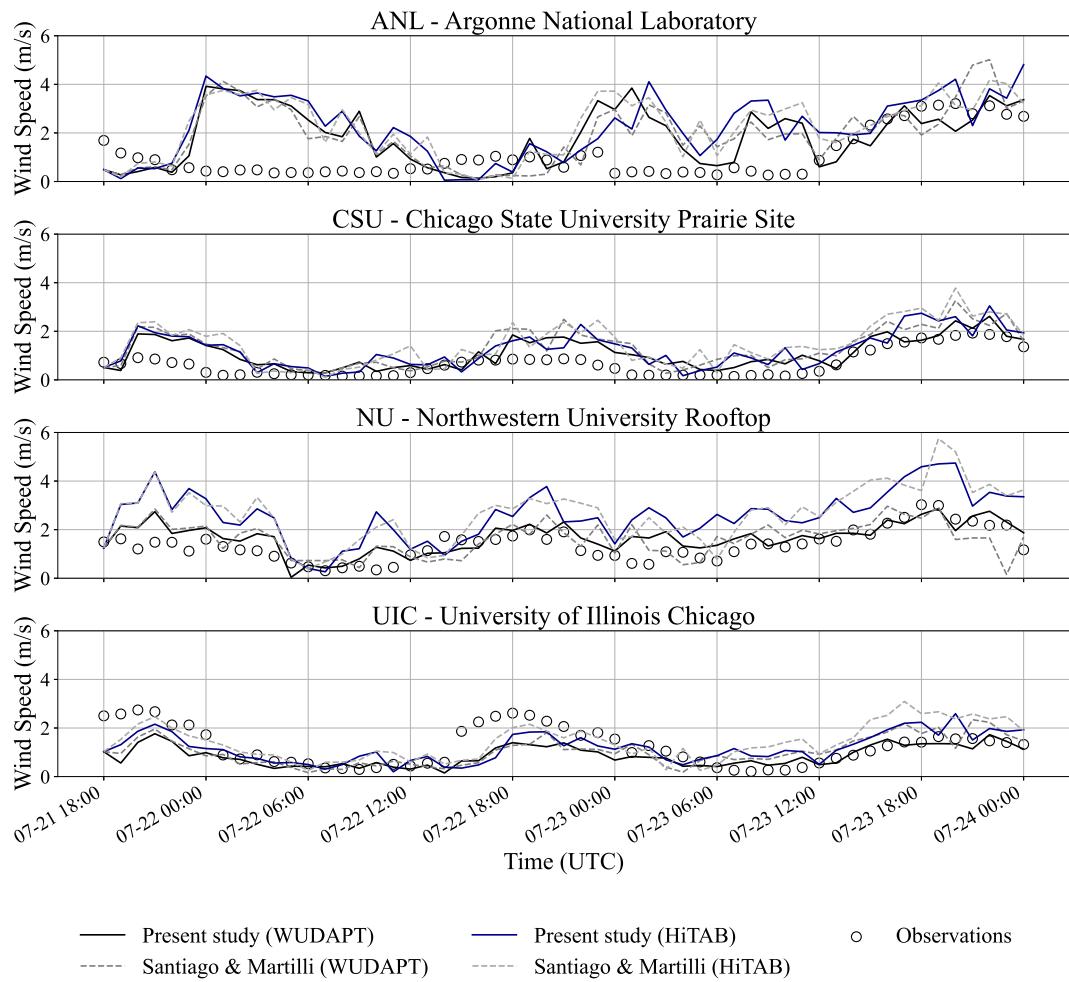
223 Figures S13 and S14 show average temperature distributions per LCZ for the WUDAPT and HiTAB experiments, respectively.  
224 In both datasets, the Present Study predicts higher mean temperatures than the Santiago and Martilli (2010) scheme in mid-  
225 density LCZs, particularly compact lowrise (LCZ3), compact midrise (LCZ2), open midrise (LCZ5), and open lowrise (LCZ6).  
226 These zones, characterized by active within-canopy turbulence, exhibit clear rightward shifts in the Present Study's PDFs and  
227 heavier right tails in the  $\Delta T$  distributions. Differences reach up to 2.4°C in WUDAPT and 1.5°C in HiTAB, highlighting the  
228 role of enhanced mixing in amplifying thermal exposure. In contrast, LCZs with sparse or industrial configurations (e.g., LCZs  
229 9, 10, and "Rest") show minimal changes, as expected due to weak or inactive canopy-layer parameterizations.

230 Figures S15 and S16 display average RH2 distributions per LCZ modeled by experiments using WUDAPT and HiTAB,  
231 respectively. The present model consistently predicts reduced RH2 in zones that also exhibit stronger warming. Compact  
232 and open midrise/lowrise LCZs show RH2 reductions of 2–6% in WUDAPT and 1–3% in HiTAB, with negatively skewed  
233  $\Delta$  RH distributions reflecting localized drying. The observed RH reductions likely reflect increased turbulent transport and  
234 mixing-driven warming near the surface, consistent with the enhanced canopy-layer turbulence captured by the updated  
235 mixing-length model. Again, differences in RH are negligible in LCZs where building-induced turbulence is relatively weaker.

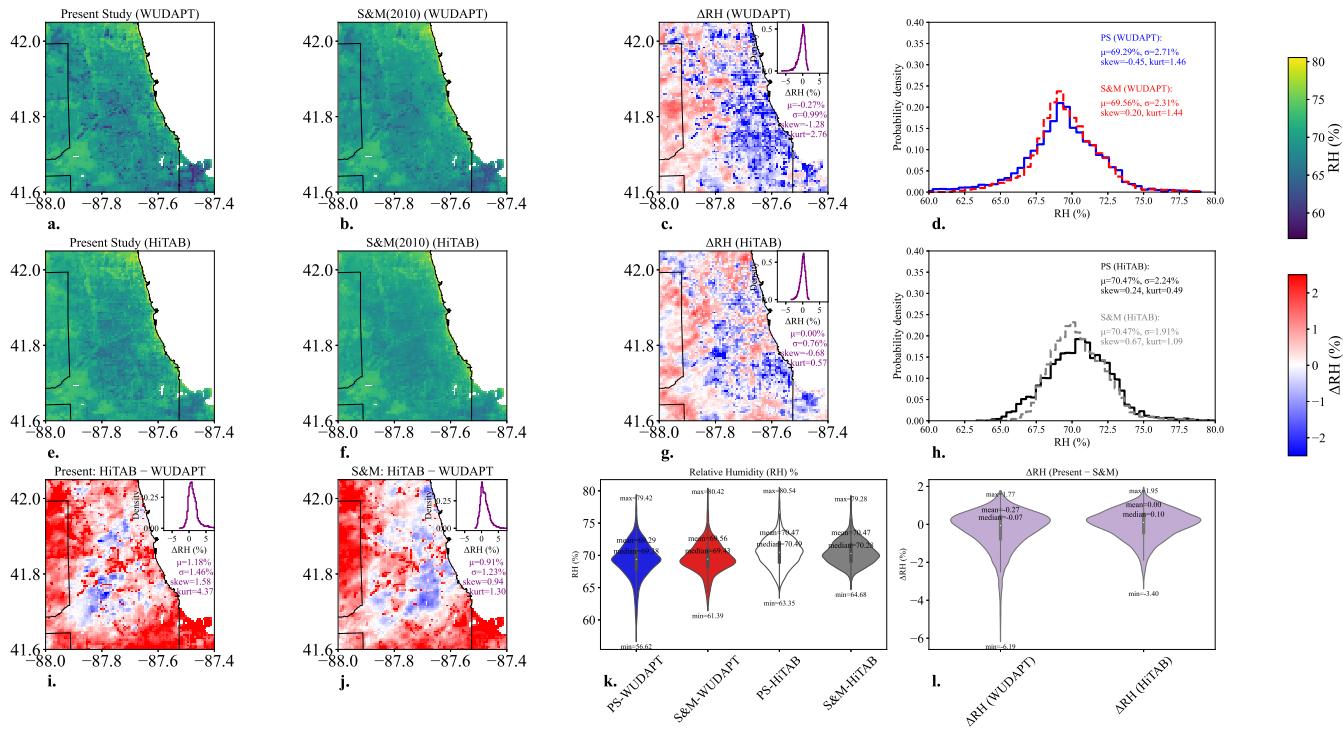
236 Figures S17 and S18 show the distribution of modeled average cooling energy demand per LCZ in the four experiments. In  
237 general, the present study yields significantly higher energy use than S&M in mid-density zones, with the strongest differences  
238 in compact and open lowrise categories (LCZs 3, 5, and 6). For WUDAPT (Fig. S17), mean per-cell  $\Delta$ Energy exceeds 2,000  
239 kWh in these categories. In HiTAB (Fig. S18), absolute values are lower due to reduced building floor counts, but the relative  
240 patterns persist, with mean differences of 500–1,000 kWh. The "Rest" category (where energy consumption is zero) and LCZs  
241 with minimal built volume show no significant differences between models.



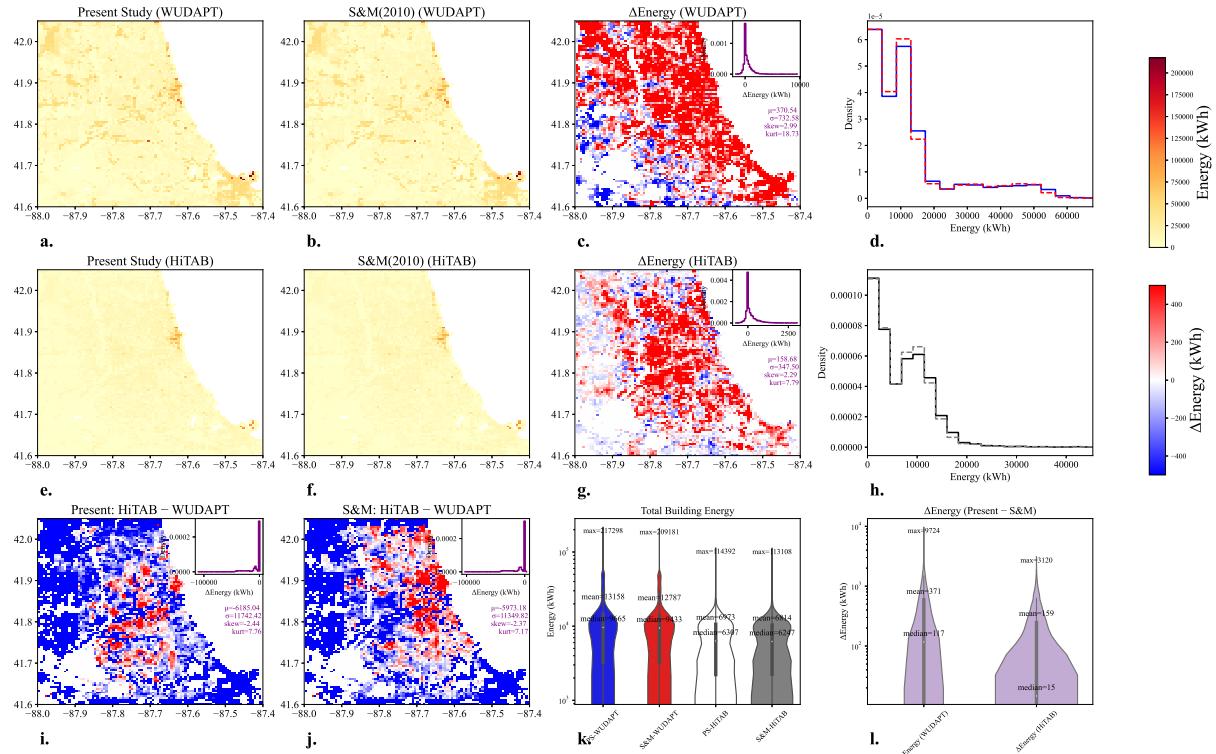
**Fig. S9.** Time series of observed and modeled near-surface air temperature at four urban observation sites (ANL, CSU, NU, and UIC) from July 21 to July 24, 2024. The present study's simulations using WUDAPT and HiTAB urban morphology datasets are shown with black and blue lines, respectively, and are compared to the Santiago and Martilli (2010) scheme (red and gray dashed lines). Observations are shown as open circles.



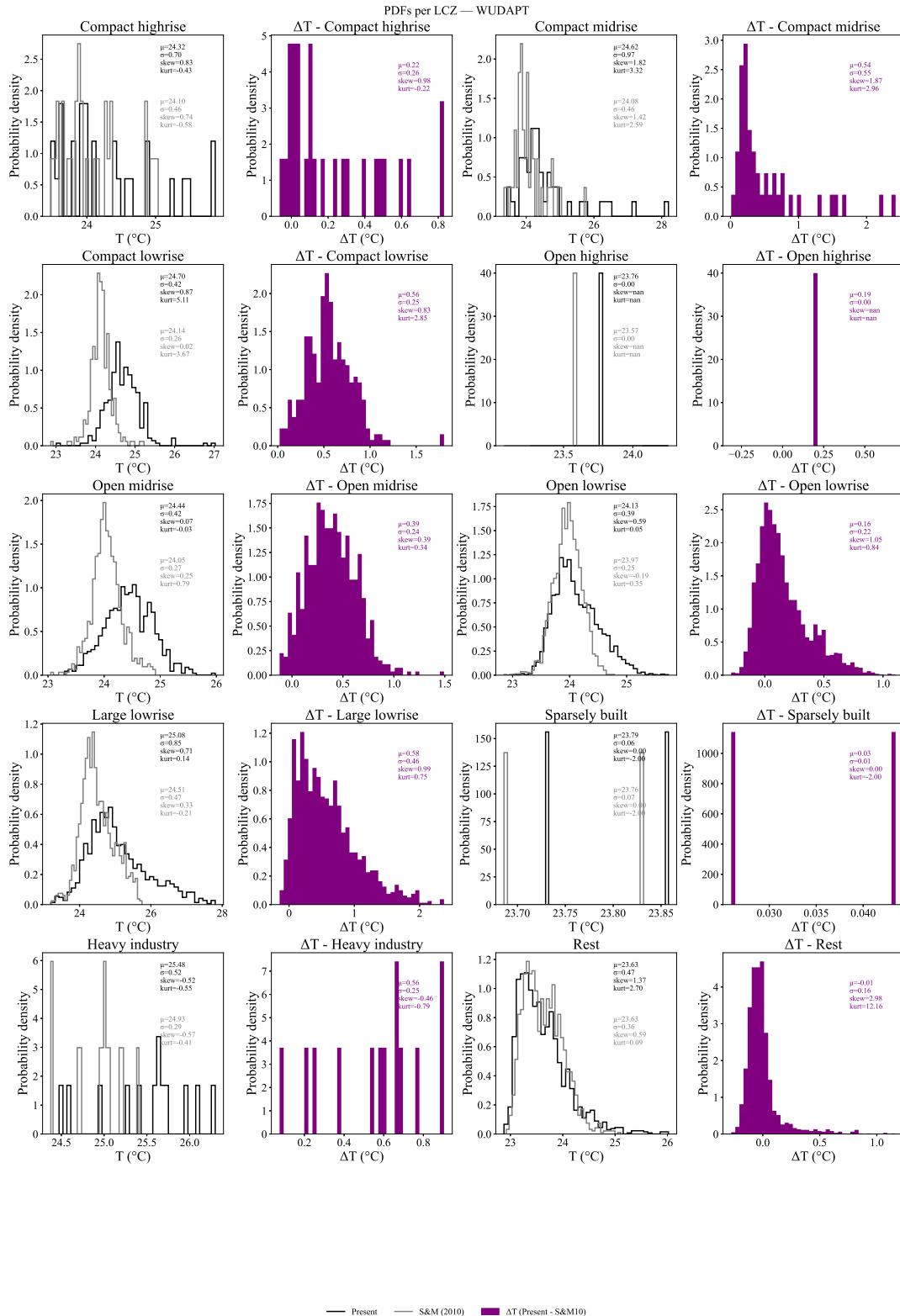
**Fig. S10.** Time series of observed and modeled wind speed at four urban observation sites (ANL, CSU, NU, and UIC) from July 21 18UTC to July 24 00UTC, 2024. The present study's simulations using WUDAPT and HiTAB urban morphology datasets are shown with black and blue lines, respectively, and are compared to the Santiago and Martilli (2010) scheme (red and gray dashed lines). Observations are shown as open circles.



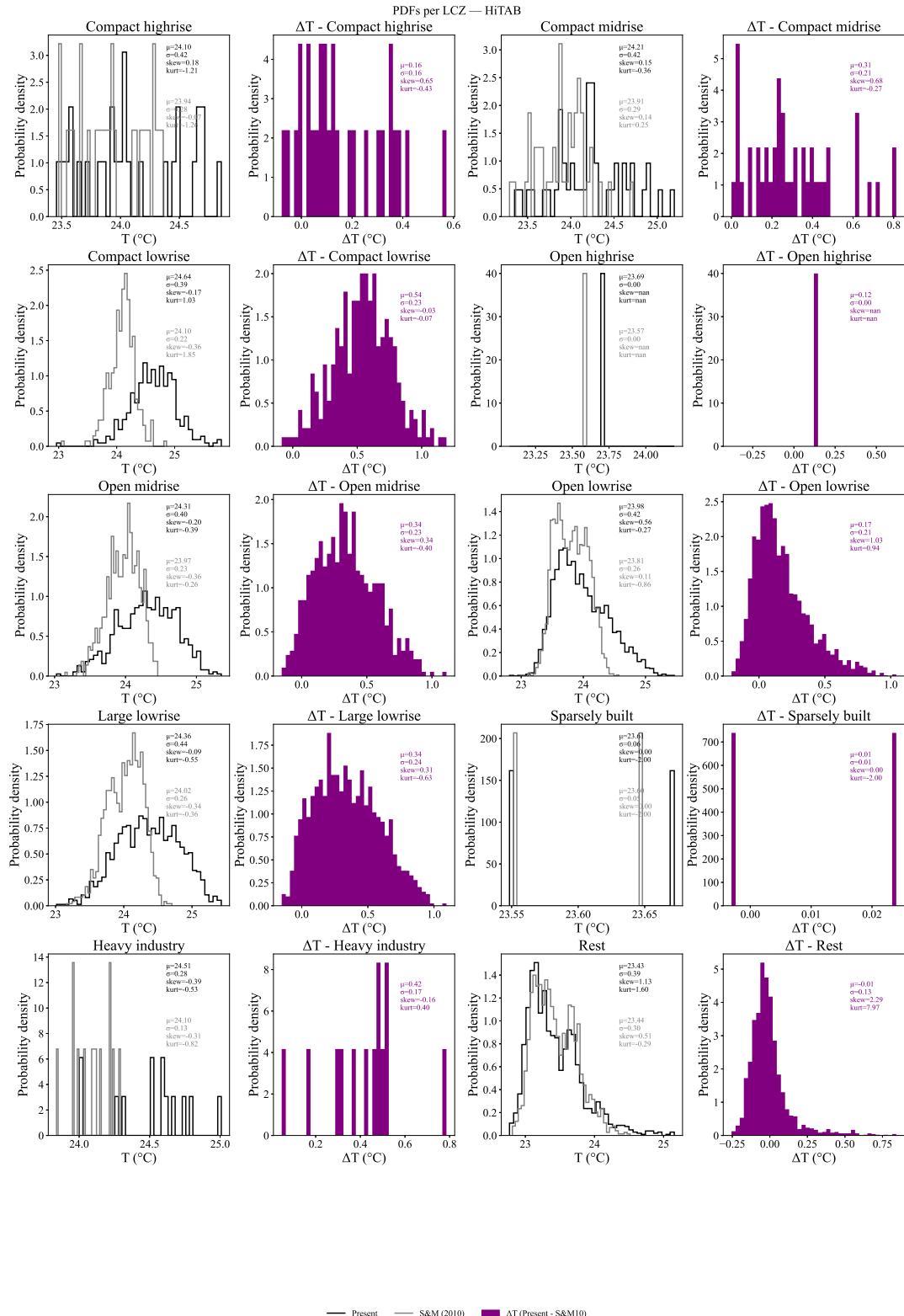
**Fig. S11.** Comparison of near-surface relative humidity (RH) from the Present Study and Santiago and Martilli 2010 (S&M) (1) using two urban morphology datasets: WUDAPT and HiTAB, over Cook County, Illinois. (a, b) Time-averaged RH fields from the Present Study and S&M using WUDAPT. (c) Spatial distribution of RH difference  $\Delta RH$  (Present – S&M) for WUDAPT, with inset showing the PDF and summary statistics. (d) Probability density functions (PDFs) of RH comparing Present Study and S&M (WUDAPT). (e, f) Same as (a, b), but using the HiTAB dataset. (g)  $\Delta RH$  (Present – S&M) for HiTAB, with inset PDF. (h) PDFs comparing Present Study and S&M using HiTAB. (i, j) Inter-dataset comparison of HiTAB – WUDAPT for Present Study and S&M, respectively, with inset PDFs and distribution metrics. (k) Violin plots of RH for all cases (Present and S&M, WUDAPT and HiTAB), showing distribution shape, mean, median, min, and max. (l) Violin plots of  $\Delta RH$  distributions between Present Study and S&M using WUDAPT and HiTAB, annotated with statistical summaries. Colorbars indicate RH (%) and  $\Delta RH$  (%).



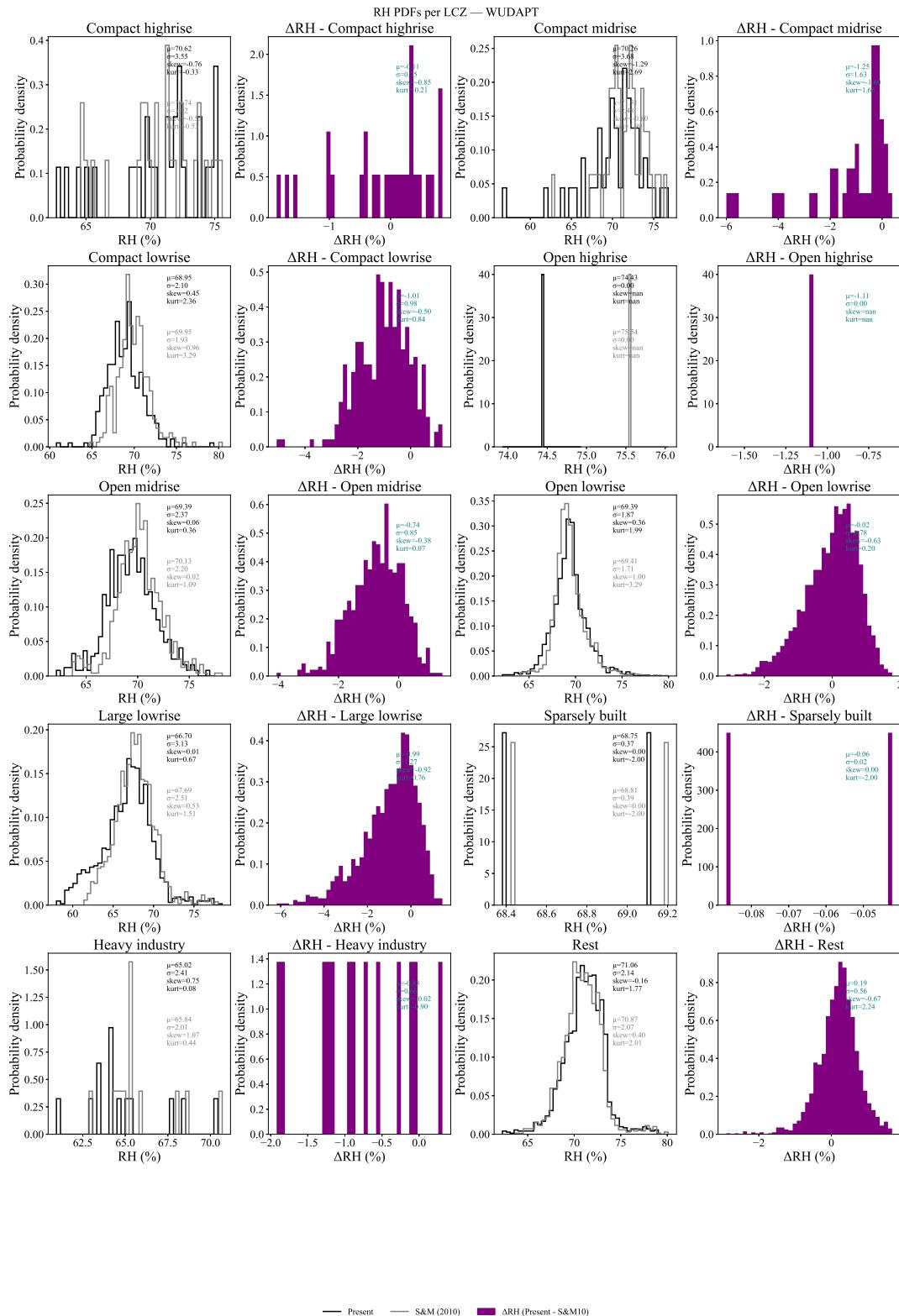
**Fig. S12.** Comparison of urban building energy consumption from the Present Study and Santiago and Martilli 2010 (S&M) (1) using two urban morphology datasets: WUDAPT and HiTAB, over Cook County, Illinois. (a, b) Total cooling energy use from the Present Study and SM using WUDAPT. (c) Spatial distribution of energy use difference  $\Delta$ Energy (Present – S&M) for WUDAPT, with inset showing the PDF and summary statistics. (d) Probability density functions (PDFs) of energy use comparing Present Study and S&M (WUDAPT). (e, f) Same as (a, b), but using the HiTAB dataset. (g)  $\Delta$ Energy (Present – S&M) for HiTAB, with inset PDF. (h) PDFs comparing Present Study and S&M using HiTAB. (i, j) Inter-dataset comparison of HiTAB – WUDAPT for Present Study and S&M, respectively, including inset PDFs and summary metrics. (k) Violin plots of total cooling energy use for all cases (Present and S&M, WUDAPT and HiTAB), showing distribution shape, mean, median, min, and max. (l) Violin plots of  $\Delta$ Energy distributions between Present Study and S&M using WUDAPT and HiTAB, annotated with statistical summaries. Colorbars indicate total energy and  $\Delta$  Energy in kilowatt-hours (kWh).



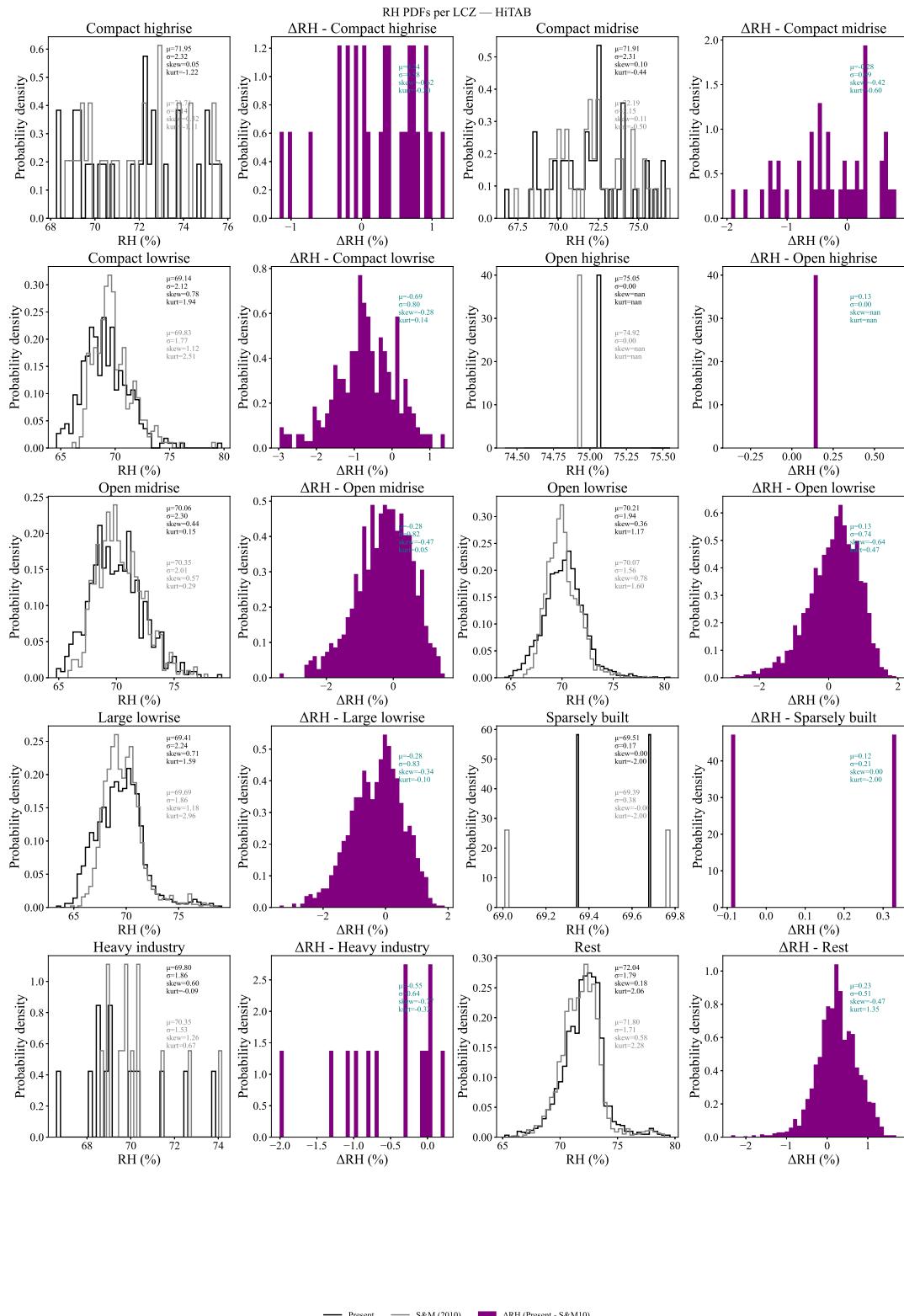
**Fig. S13.** Probability density functions (PDFs) of near-surface temperature ( $T$ ) for different Local Climate Zones (LCZs), comparing the present study (black lines) and Santiago and Martilli (2010) (gray lines). Each row represents different LCZ categories. Middle panels show PDFs of the temperature differences ( $\Delta T$ ) between the two modeling approaches (present study minus Santiago and Martilli, 2010). Statistics displayed include mean ( $\mu$ ), standard deviation ( $\sigma$ ), skewness (skew), and kurtosis (kurt).



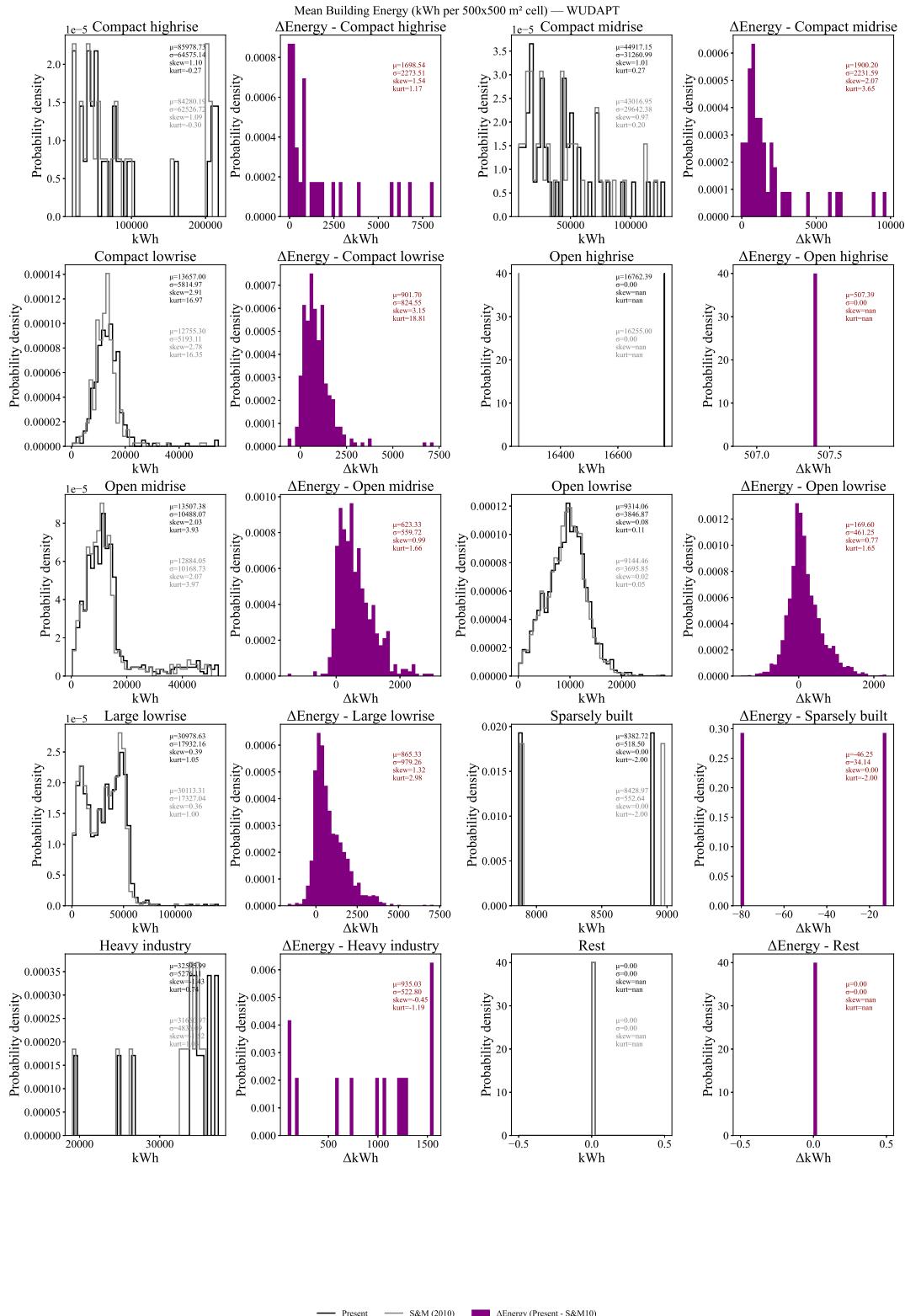
**Fig. S14.** Probability density functions (PDFs) of near-surface temperature ( $T$ ) for different Local Climate Zones (LCZs), comparing the present study (black lines) and Santiago and Martilli (2010) (gray lines). Each row represents different LCZ categories. Middle panels show PDFs of the temperature differences ( $\Delta T$ ) between the two modeling approaches (present study minus Santiago and Martilli, 2010). Statistics displayed include mean ( $\mu$ ), standard deviation ( $\sigma$ ), skewness (skew), and kurtosis (kurt).



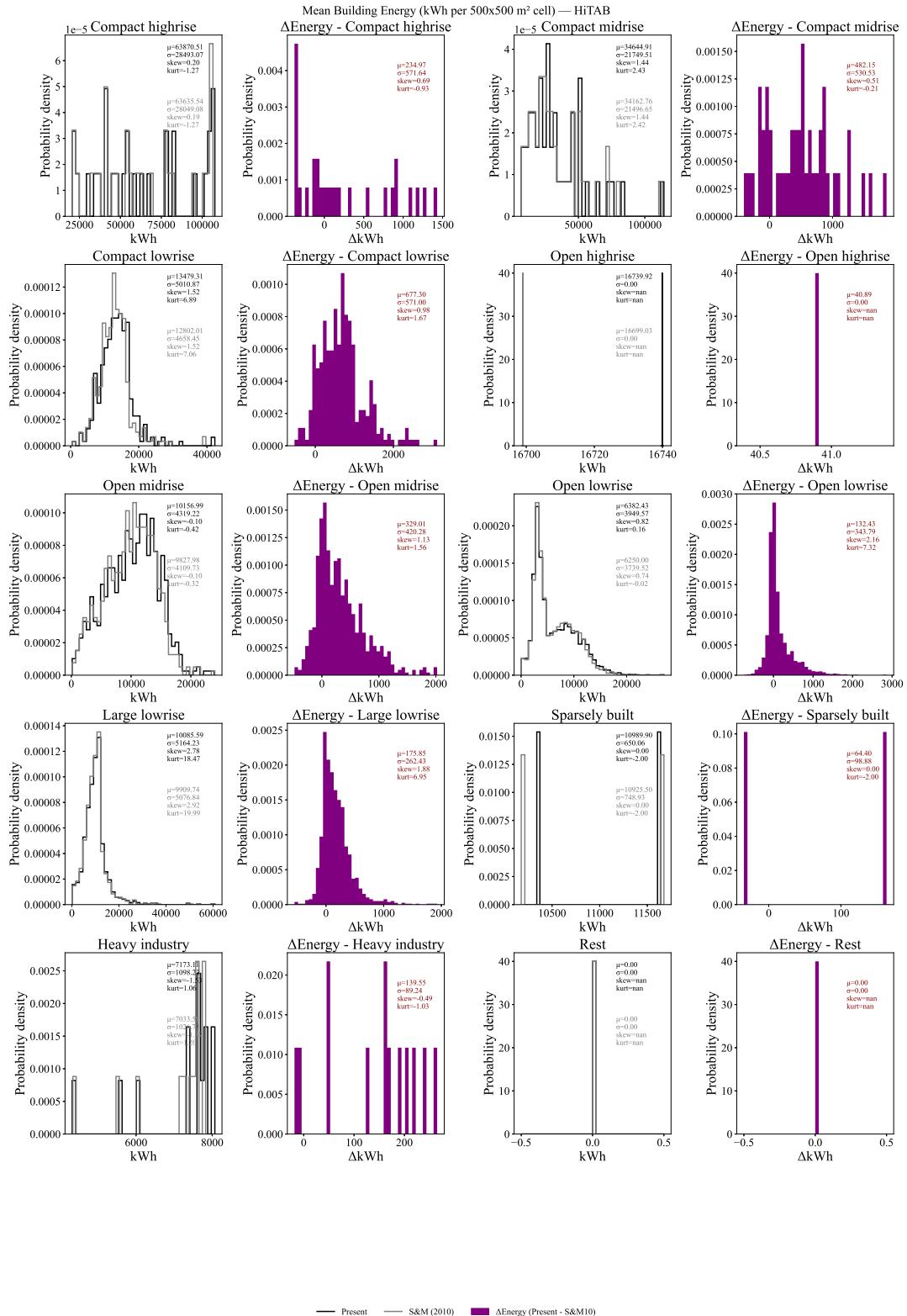
**Fig. S15.** Probability density functions (PDFs) of near-surface temperature ( $T$ ) for different Local Climate Zones (LCZs), comparing the present study (black lines) and Santiago and Martilli (2010) (gray lines). Each row represents different LCZ categories. Middle panels show PDFs of the temperature differences ( $\Delta T$ ) between the two modeling approaches (present study minus Santiago and Martilli, 2010). Statistics displayed include mean ( $\mu$ ), standard deviation ( $\sigma$ ), skewness (skew), and kurtosis (kurt).



**Fig. S16.** Probability density functions (PDFs) of near-surface temperature ( $T$ ) for different Local Climate Zones (LCZs), comparing the present study (black lines) and Santiago and Martilli (2010) (gray lines). Each row represents different LCZ categories. Middle panels show PDFs of the temperature differences ( $\Delta T$ ) between the two modeling approaches (present study minus Santiago and Martilli, 2010). Statistics displayed include mean ( $\mu$ ), standard deviation ( $\sigma$ ), skewness (skew), and kurtosis (kurt).



**Fig. S17.** Probability density functions (PDFs) of near-surface temperature ( $T$ ) for different Local Climate Zones (LCZs), comparing the present study (black lines) and Santiago and Martilli (2010) (gray lines). Each row represents different LCZ categories. Middle panels show PDFs of the temperature differences ( $\Delta T$ ) between the two modeling approaches (present study minus Santiago and Martilli, 2010). Statistics displayed include mean ( $\mu$ ), standard deviation ( $\sigma$ ), skewness (skew), and kurtosis (kurt).



**Fig. S18.** Probability density functions (PDFs) of near-surface temperature ( $T$ ) for different Local Climate Zones (LCZs), comparing the present study (black lines) and Santiago and Martilli (2010) (gray lines). Each row represents different LCZ categories. Middle panels show PDFs of the temperature differences ( $\Delta T$ ) between the two modeling approaches (present study minus Santiago and Martilli, 2010). Statistics displayed include mean ( $\mu$ ), standard deviation ( $\sigma$ ), skewness (skew), and kurtosis (kurt).

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