

1      **Non-line-of-sight imaging of moving objects obscured by a**  
 2      **random corridor**

3      **Supplementary Information**

4      **Supplementary Note 1:**

5      **Imaging principle of RSESC method**

6      In the paper, the cross-correlation of speckle intensity images  $C_{rd}(\Delta\mathbf{p})$   
 7      calculated by the RSESC method is approximately equivalent to the transmitted field  
 8      autocorrelation of a hidden object  $C_{ho}(\Delta\mathbf{p})$  under the condition of multiple scattering.  
 9      Hidden objects in random corridors can be imaged with the approximation. The  
 10     approximate relationship can be expressed as,

11     
$$C_{rd}(\Delta\mathbf{p}) \approx C_{ho}(\Delta\mathbf{p}) \quad (1)$$

12     This relationship can be derived from the definition of  $C_{rd}(\Delta\mathbf{p})$  based on three  
 13     assumptions.  $C_{rd}(\Delta\mathbf{p})$  is defined as,

14     
$$C_{rd}(\Delta\mathbf{r}) = \left\langle \frac{\left[ I_{rd}(\mathbf{r}_c; \mathbf{r}_o) - \langle I_{rd}(\mathbf{r}_c; \mathbf{r}_o) \rangle_{\mathbf{r}_c} \right] \left[ I_{rd}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) - \langle I_{rd}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \rangle_{\mathbf{r}_c} \right]}{\sigma_{rd}(\mathbf{r}_o) \sigma_{rd}(\mathbf{r}_o + \Delta\mathbf{p})} \right\rangle_{\mathbf{r}_c, \mathbf{r}_o} \quad (2)$$

15     The ensemble average  $\langle \dots \rangle_{\mathbf{r}_o}$  is used to increase the overall average times. Therefore,  
 16      $C_{rd}(\Delta\mathbf{p})$  is expressed as,

17     
$$C_{rd}(\Delta\mathbf{r}) = \frac{\left[ I_{rd}(\mathbf{r}_c; \mathbf{r}_o) - \langle I_{rd}(\mathbf{r}_c; \mathbf{r}_o) \rangle_{\mathbf{r}_c} \right] \left[ I_{rd}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) - \langle I_{rd}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \rangle_{\mathbf{r}_c} \right]}{\sigma_{rd}(\mathbf{r}_o) \sigma_{rd}(\mathbf{r}_o + \Delta\mathbf{p})} \quad (3)$$

18     The numerator is expanded in formula (3) to obtain,

19     
$$C_{rd}(\Delta\mathbf{r}) = \frac{\langle I_{rd}(\mathbf{r}_c; \mathbf{r}_o) I_{rd}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \rangle_{\mathbf{r}_c} - \langle I_{rd}(\mathbf{r}_c; \mathbf{r}_o) \rangle_{\mathbf{r}_c} \langle I_{rd}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \rangle_{\mathbf{r}_c}}{\sigma_{rd}(\mathbf{r}_o) \sigma_{rd}(\mathbf{r}_o + \Delta\mathbf{p})} \quad (4)$$

20     The reduced electric field of scattered light is denoted as  $E_{rd}$ , and then the reduced  
 21     intensity of scattered light can be expanded as  $I_{rd} = E_{rd} E_{rd}^*$ .

22 **Assumption 1:**  $E_{rd}$  follows a Gaussian distribution with zero mean. Apply a moment  
 23 theorem for complex Gaussian<sup>1</sup>, and the first term of the numerator in Eq(4) can be  
 24 expanded as,

$$\begin{aligned}
 & \left\langle I_{rd}(\mathbf{r}_c; \mathbf{r}_o) I_{rd}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} \\
 &= \left\langle E_{rd}(\mathbf{r}_c; \mathbf{r}_o) E_{rd}^*(\mathbf{r}_c; \mathbf{r}_o) E_{rd}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) E_{rd}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} \\
 25 &= \left| \left\langle E_{rd}(\mathbf{r}_c; \mathbf{r}_o) E_{rd}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} \right|^2 + \left\langle \left| E_{rd}(\mathbf{r}_c; \mathbf{r}_o) \right|^2 \right\rangle_{\mathbf{r}_c} \left\langle \left| E_{rd}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right|^2 \right\rangle_{\mathbf{r}_c} \quad (5) \\
 &= \left| \left\langle E_{rd}(\mathbf{r}_c; \mathbf{r}_o) E_{rd}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} \right|^2 + \left\langle I_{rd}(\mathbf{r}_c; \mathbf{r}_o) \right\rangle_{\mathbf{r}_c} \left\langle I_{rd}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c}
 \end{aligned}$$

26 At the same time, the denominator in Eq. (4) reads<sup>2</sup>,

$$\begin{aligned}
 & \sigma_{rd}(\mathbf{r}_o) \sigma_{rd}(\mathbf{r}_o + \Delta\mathbf{p}) \\
 27 &= \left\langle I_{rd}(\mathbf{r}_c; \mathbf{r}_o) \right\rangle_{\mathbf{r}_c} \left\langle I_{rd}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} \quad (6) \\
 &= \left\langle \left| E_{rd}(\mathbf{r}_c; \mathbf{r}_o) \right|^2 \right\rangle_{\mathbf{r}_c} \left\langle \left| E_{rd}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right|^2 \right\rangle_{\mathbf{r}_c}
 \end{aligned}$$

28 Substituting (5) and (6) into (4), we have,

$$29 C_{rd}(\Delta\mathbf{r}) = \frac{\left\langle E_{rd}(\mathbf{r}_c; \mathbf{r}_o) E_{rd}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c}}{\left\langle \left| E_{rd}(\mathbf{r}_c; \mathbf{r}_o) \right|^2 \right\rangle_{\mathbf{r}_c}} \cdot \frac{\left\langle E_{rd}^*(\mathbf{r}_c; \mathbf{r}_o) E_{rd}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c}}{\left\langle \left| E_{rd}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right|^2 \right\rangle_{\mathbf{r}_c}} \quad (7)$$

30 We denote the transmitted field of the hidden object as  $E_{ho}$ . Previous studies have  
 31 shown that  $E_{rd}$  is a function of  $E_{ho}$ <sup>3</sup>,

$$32 E_{rd}(\mathbf{r}_c; \mathbf{r}_o) = \int d\mathbf{r} K_{rd}(\mathbf{r}_c, \mathbf{r}) \cdot E_{ho}(\mathbf{r}; \mathbf{r}_o) \quad (8)$$

$$33 E_{rd}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) = \int d\mathbf{r} K_{rd}(\mathbf{r}_c, \mathbf{r}) \cdot E_{ho}(\mathbf{r}; \mathbf{r}_o + \Delta\mathbf{p}) \quad (9)$$

34 Here  $\mathbf{r}$  represents the position vector in three-dimensional space with random media,  
 35  $K_{rd}(\mathbf{r}_c, \mathbf{r})$  represents the subspace reduced generalized propagator (SRGP) from  $\mathbf{r}$  to  
 36  $\mathbf{r}_c$ . Substituting (8) and (9) into (7),

$$\begin{aligned}
 & \left\langle E_{rd}(\mathbf{r}_c; \mathbf{r}_o) E_{rd}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} \\
 37 &= \left\langle \int E_{ho}(\mathbf{r}'; \mathbf{r}_o) K_{rd}(\mathbf{r}_c, \mathbf{r}') d\mathbf{r}' \int E_{ho}^*(\mathbf{r}''; \mathbf{r}_o + \Delta\mathbf{p}) K_{rd}^*(\mathbf{r}_c, \mathbf{r}'') d\mathbf{r}'' \right\rangle_{\mathbf{r}_c} \quad (10) \\
 &= \left\langle \int \int E_{ho}(\mathbf{r}'; \mathbf{r}_o) E_{ho}^*(\mathbf{r}''; \mathbf{r}_o + \Delta\mathbf{p}) K_{rd}(\mathbf{r}_c, \mathbf{r}') K_{rd}^*(\mathbf{r}_c, \mathbf{r}'') d\mathbf{r}' d\mathbf{r}'' \right\rangle_{\mathbf{r}_c}
 \end{aligned}$$

38 **Assumption 2:** When  $\mathbf{r}' \neq \mathbf{r}''$ ,  $\int K_{\text{rd}}(\mathbf{r}_c, \mathbf{r}') K_{\text{rd}}^*(\mathbf{r}_c, \mathbf{r}'') d\mathbf{r}' = 0$ , namely, SRGP from  
 39 different spatial locations to  $\mathbf{r}_c$  is independent of each other. Therefore, equation (10)  
 40 can be reduced to,

$$\begin{aligned}
 & \left\langle E_{\text{rd}}(\mathbf{r}_c; \mathbf{r}_o) E_{\text{rd}}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} \\
 41 &= \left\langle \int E_{\text{ho}}(\mathbf{r}; \mathbf{r}_o) E_{\text{ho}}^*(\mathbf{r}; \mathbf{r}_o + \Delta\mathbf{p}) |K_{\text{rd}}(\mathbf{r}_c, \mathbf{r})|^2 d\mathbf{r} \right\rangle_{\mathbf{r}_c} \\
 &= \int E_{\text{ho}}(\mathbf{r}; \mathbf{r}_o) E_{\text{ho}}^*(\mathbf{r}; \mathbf{r}_o + \Delta\mathbf{p}) \left\langle |K_{\text{rd}}(\mathbf{r}_c, \mathbf{r})|^2 \right\rangle_{\mathbf{r}_c} d\mathbf{r}
 \end{aligned} \tag{11}$$

42 **Assumption 3:**  $P_G = \left\langle |K_{\text{rd}}(\mathbf{r}_c, \mathbf{r})|^2 \right\rangle_{\mathbf{r}_c}$  is a constant and independent of  $\mathbf{r}$  and  $\mathbf{r}_c$ , then,

$$43 \quad \left\langle E_{\text{rd}}(\mathbf{r}_c; \mathbf{r}_o) E_{\text{rd}}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} = P_G \int E_{\text{ho}}(\mathbf{r}; \mathbf{r}_o) E_{\text{ho}}^*(\mathbf{r}; \mathbf{r}_o + \Delta\mathbf{p}) d\mathbf{r} \tag{12}$$

44 Similarly,

$$45 \quad \left\langle |E_{\text{rd}}(\mathbf{r}_c; \mathbf{r}_o)|^2 \right\rangle_{\mathbf{r}_c} = P_G \int |E_{\text{ho}}(\mathbf{r}; \mathbf{r}_o)|^2 d\mathbf{r} \tag{13}$$

$$46 \quad \left\langle |E_{\text{rd}}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p})|^2 \right\rangle_{\mathbf{r}_c} = P_G \int |E_{\text{ho}}(\mathbf{r}; \mathbf{r}_o + \Delta\mathbf{p})|^2 d\mathbf{r} \tag{14}$$

47 Therefore, it can be approximated as,

$$48 \quad \frac{\left\langle E_{\text{rd}}(\mathbf{r}_c; \mathbf{r}_o) E_{\text{rd}}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c}}{\left\langle |E_{\text{rd}}(\mathbf{r}_c; \mathbf{r}_o)|^2 \right\rangle_{\mathbf{r}_c}} \approx \frac{\int E_{\text{ho}}(\mathbf{r}; \mathbf{r}_o) E_{\text{ho}}^*(\mathbf{r}; \mathbf{r}_o + \Delta\mathbf{p}) d\mathbf{r}}{\int |E_{\text{ho}}(\mathbf{r}; \mathbf{r}_o)|^2 d\mathbf{r}} \tag{15}$$

49 We define the autocorrelation of the transmitted field of a hidden object as,

$$50 \quad C_{\text{ho}}(\Delta\mathbf{p}) = \left| \frac{\int E_{\text{ho}}(\mathbf{r}; \mathbf{r}_o) E_{\text{ho}}^*(\mathbf{r}; \mathbf{r}_o + \Delta\mathbf{p}) d\mathbf{r}}{\int |E_{\text{ho}}(\mathbf{r}; \mathbf{r}_o)|^2 d\mathbf{r}} \right|^2 \tag{16}$$

51 By substituting Eq. (15), (16) into Eq. (7), eventually we obtain the imaging principle  
 52 of the RSEVIC method,

$$53 \quad C_{\text{rd}}(\Delta\mathbf{r}) \approx C_{\text{ho}}(\Delta\mathbf{p}) = \left| \frac{\int E_{\text{ho}}(\mathbf{r}; \mathbf{r}_o) E_{\text{ho}}^*(\mathbf{r}; \mathbf{r}_o + \Delta\mathbf{p}) d\mathbf{r}}{\int |E_{\text{ho}}(\mathbf{r}; \mathbf{r}_o)|^2 d\mathbf{r}} \right|^2 \tag{17}$$

54 Assumption 2 is satisfied by strong scattering of random media. The subspace  
 55 reduction we proposed guarantees that Assumption 3 is satisfied regardless of the shape  
 56 of random medium. Moreover, subspace reduction also ensures the validity of  
 57 Assumption 1 for any observation position  $\mathbf{r}_c$  and any size of spatial statistical region  
 58  $\langle \dots \rangle_{\mathbf{r}_c}$ .

59 In expression (16),  $\int d\mathbf{r}$  represents the integral of three-dimensional space  
 60 including random medium,  $E_{ho}(\mathbf{r}; \mathbf{r}_o)$  and  $E_{ho}(\mathbf{r}; \mathbf{r}_o + \Delta\mathbf{p})$  represent the three-  
 61 dimensional transmitted field illuminating on a hidden object. Therefore,

$$62 \quad \begin{aligned} & \int E_{ho}(\mathbf{r}; \mathbf{r}_o) E_{ho}^*(\mathbf{r}; \mathbf{r}_o + \Delta\mathbf{p}) d\mathbf{r} \\ &= \iiint E_{ho}(x, y, z; \mathbf{r}_o) E_{ho}^*(x, y, z; \mathbf{r}_o + \Delta\mathbf{p}) dx dy dz \end{aligned} \quad (18)$$

63 When the hidden object is a transmission plate and the integral interval of  $\int dz$  is  
 64 relatively small,  $E_{ho}(x, y, z) \approx E_{ho}(x, y)$ , thus,

$$65 \quad \begin{aligned} & \int E_{ho}(\mathbf{r}; \mathbf{r}_o) E_{ho}^*(\mathbf{r}; \mathbf{r}_o + \Delta\mathbf{p}) d\mathbf{r} \\ &= b_z \cdot \iint E_{ho}(x, y; \mathbf{r}_o) E_{ho}^*(x, y; \mathbf{r}_o + \Delta\mathbf{p}) dx dy \end{aligned} \quad (19)$$

66 Here  $b_z = \int dz$  is a constant. Applying the Eq. (19) to the Eq. (16), we have,

$$67 \quad C_{ho}(\Delta\mathbf{p}) = \left| \frac{\iint E_{ho}(x, y; \mathbf{r}_o) E_{ho}^*(x, y; \mathbf{r}_o + \Delta\mathbf{p}) dx dy}{\iint |E_{ho}(x, y; \mathbf{r}_o)|^2 dx dy} \right|^2 \quad (20)$$

68 We define the reduced dimensionless transmitted field as,

$$69 \quad \hat{E}_{ho}(x, y; \mathbf{r}_o) = \frac{E_{ho}(x, y; \mathbf{r}_o)}{\sqrt{\iint |E_{ho}(x, y; \mathbf{r}_o)|^2 dx dy}} \quad (21)$$

$$70 \quad \hat{E}_{ho}(x, y; \mathbf{r}_o + \Delta\mathbf{p}) = \frac{E_{ho}(x, y; \mathbf{r}_o + \Delta\mathbf{p})}{\sqrt{\iint |E_{ho}(x, y; \mathbf{r}_o + \Delta\mathbf{p})|^2 dx dy}} \quad (22)$$

71 By inserting Eq. (21), (22) to Eq. (20) and applying the autocorrelation theorem<sup>4</sup>, we  
 72 have,

$$\begin{aligned}
C_{\text{ho}}(\Delta\mathbf{p}) &= \left| \iint \hat{E}_{\text{ho}}(x, y; \mathbf{r}_o) \hat{E}_{\text{ho}}^*(x, y; \mathbf{r}_o + \Delta\mathbf{p}) dx dy \right|^2 \\
73 \quad &= \left| \mathcal{F}^{-1} \left\{ \left| \mathcal{F} \left\{ \hat{E}_{\text{ho}}(x, y) \right\} \right|^2 \right\} \right|^2
\end{aligned} \tag{23}$$

74 Here  $\mathcal{F}$  and  $\mathcal{F}^{-1}$  represent Fourier transform and inverse Fourier transform,  
75 respectively. Then combining Eq. (23) with (1),

$$76 \quad \hat{E}_{\text{ho}}(x, y) = \mathcal{F}^{-1} \left\{ \sqrt{\mathcal{F} \left\{ \sqrt{C_{\text{rd}}(\Delta\mathbf{p})} \cdot e^{i\varphi_1(\Delta\mathbf{p})} \right\}} \cdot e^{i\varphi_2(\mathbf{k})} \right\} \tag{24}$$

77 In this equation,  $\varphi_1(\Delta\mathbf{p})$  is the phase of  $\iint \hat{E}_{\text{ho}}(x, y; \mathbf{r}_o) \hat{E}_{\text{ho}}^*(x, y; \mathbf{r}_o + \Delta\mathbf{p}) dx dy$ . When  
78 the hidden object is a transmission plate, we find that the imaginary part of  
79  $\iint \hat{E}_{\text{ho}}(x, y; \mathbf{r}_o) \hat{E}_{\text{ho}}^*(x, y; \mathbf{r}_o + \Delta\mathbf{p}) dx dy$  is much smaller than its real part, so  
80  $\varphi_1(\Delta\mathbf{p}) \approx 0$ .  $\varphi_2(\mathbf{k})$  can be recovered by an iterative phase recovery algorithm<sup>5, 6</sup>.  
81 Substituting  $\varphi_1(\Delta\mathbf{p})$ ,  $\varphi_2(\mathbf{k})$  and  $C_{\text{rd}}(\Delta\mathbf{p})$  into (24),  $\hat{E}_{\text{ho}}(x, y)$  can be recovered.  
82  $\hat{E}_{\text{ho}}(x, y)$  is the reduced dimensionless transmitted field on an XY-plane near the  
83 hidden object, and its amplitude  $|\hat{E}_{\text{ho}}(x, y)|$  is the image of the hidden object. Finally,  
84 the image of the hidden object is successfully reconstructed.

85 **Supplementary Note 2:**

86 **Calculation of the autocorrelation of hidden objects**

87 In our experiment, the hidden object to be imaged is a transmission plate. The  
88 transmittance of the aperture of the transmission plate to the laser is 1, and the  
89 transmittance of other parts of the plate is close to 0. Because the size of the aperture  
90 on the transmission plate is in the order of mm, which is  $10^3$  times of the laser  
91 wavelength, the diffraction effect can be ignored. Based on the analysis above, the  
92 transmitted field of different hidden objects can be simulated.

93 The steps to simulate the transmitted field of a hidden object are as follows. First,  
94 an amplitude function  $A_{\text{ho}}(x, y)$  is generated. Then the center of the transmission  
95 plate is taken as the coordinate origin, and  $A_{\text{ho}}(x, y)$  is binarized according to the

96 geometric characteristics of the transmission plate. The rule of binarization is,

97

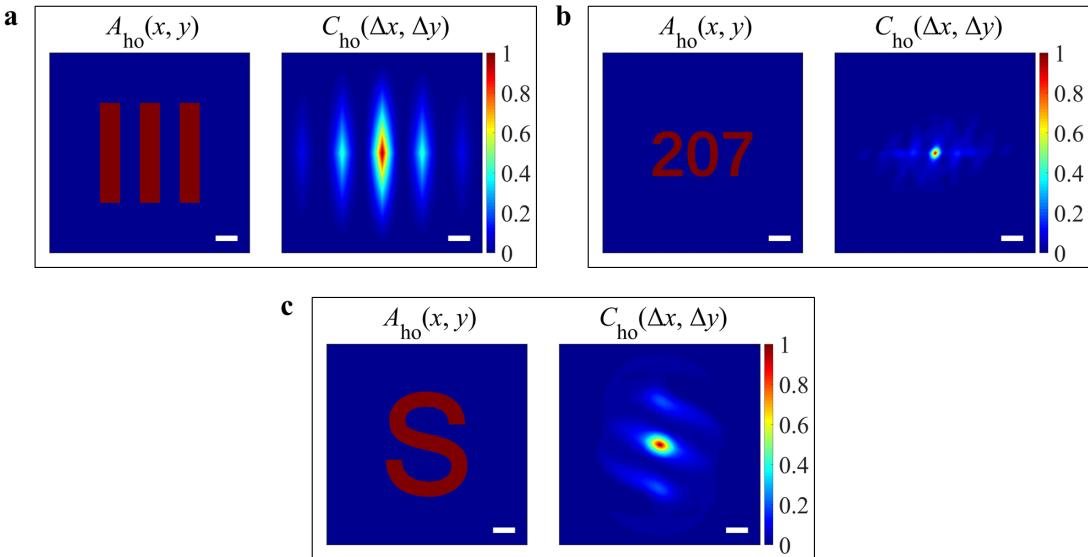
$$A_{\text{ho}}(x, y) = \begin{cases} 1, & (x, y) \in \text{Hole area} \\ 0, & (x, y) \notin \text{Hole area} \end{cases} \quad (25)$$

98 Next, the phase of the transmitted field of the hidden object is simulated as a constant  
 99  $\varphi_{\text{ho}}$  that is independent of  $x$  or  $y$ . Finally, the simulated transmitted field of the  
 100 hidden object is obtained,

101

$$E_{\text{ho}}(x, y) = A_{\text{ho}}(x, y) \cdot e^{i\varphi_{\text{ho}}} \quad (26)$$

102 Substituting the simulated  $E_{\text{ho}}(x, y)$  into Eq. (21), (23),  $C_{\text{ho}}(\Delta x, \Delta y)$  can be  
 103 obtained. The simulated  $A_{\text{ho}}(x, y)$  and  $C_{\text{ho}}(\Delta x, \Delta y)$  of three different hidden objects  
 104 are shown in the supplementary Figure 1.



106  
 107 Supplementary Figure 1: Simulated  $A_{\text{ho}}(x, y)$  and  $C_{\text{ho}}(\Delta x, \Delta y)$  of three different hidden  
 108 objects. Scale bar: 1mm.

109 **Supplementary Note 3:**

110 **Elimination of the adverse effect of ambient noise**

111 Since the Eq.1 of the imaging relation is obtained without concerns for the noise,  
 112 it is no longer valid when the speckle intensity images contain noise. In the condition of  
 113 taking the noise into consideration, the reduced electric field  $E_{\text{rd}}$  of scattered light is  
 114 expressed as,

115

$$E_{\text{rd}}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) = E_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) + E_{\text{rd-b}}(\mathbf{r}_c; t_0 + \Delta t) \quad (27)$$

116 Here  $E_{\text{rd-ho}}$  represents the scattering field induced by the transmitted field  $E_{\text{ho}}$   
 117 interacting with the random medium.  $E_{\text{rd-b}}$  represents the time-varying background  
 118 noise field. The position vector of the hidden object at time  $t_0$  is  $\mathbf{r}_o$ , and the velocity  
 119 vector of the hidden object is  $\mathbf{v}$ , then  $\Delta\mathbf{p}=\mathbf{v}\Delta t$ . Based on equation (27), the cross-  
 120 correlation of  $E_{\text{rd}}$  can be expanded into four terms,

$$\begin{aligned} & \left\langle E_{\text{rd}}(\mathbf{r}_c; \mathbf{r}_o) E_{\text{rd}}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} \\ 121 &= \left\langle E_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) E_{\text{rd-ho}}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) + E_{\text{rd-b}}(\mathbf{r}_c; t_0) E_{\text{rd-b}}^*(\mathbf{r}_c; t_0 + \Delta t) + \right. \\ & \quad \left. E_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) E_{\text{rd-b}}^*(\mathbf{r}_c; t_0 + \Delta t) + E_{\text{rd-b}}(\mathbf{r}_c; t_0) E_{\text{rd-ho}}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} \end{aligned} \quad (28)$$

122 Since  $E_{\text{rd-ho}}$  and  $E_{\text{rd-b}}$  are independent of each other,

$$123 \quad \left\langle E_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) E_{\text{rd-b}}^*(\mathbf{r}_c; t_0 + \Delta t) + E_{\text{rd-b}}(\mathbf{r}_c; t_0) E_{\text{rd-ho}}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} = 0 \quad (29)$$

124 Substituting (29) into (28), then

$$\begin{aligned} & \left\langle E_{\text{rd}}(\mathbf{r}_c; \mathbf{r}_o) E_{\text{rd}}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle \\ 125 &= \left\langle E_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) E_{\text{rd-ho}}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} + \left\langle E_{\text{rd-b}}(\mathbf{r}_c; t_0) E_{\text{rd-b}}^*(\mathbf{r}_c; t_0 + \Delta t) \right\rangle_{\mathbf{r}_c} \end{aligned} \quad (30)$$

126 Then substituting (30) into (7), we have

$$\begin{aligned} C_{\text{rd}}(\Delta\mathbf{p}) &= \frac{\left| \left\langle E_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) E_{\text{rd-ho}}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} \right|^2 + \left| \left\langle E_{\text{rd-b}}(\mathbf{r}_c; t_0) E_{\text{rd-b}}^*(\mathbf{r}_c; t_0 + \Delta t) \right\rangle_{\mathbf{r}_c} \right|^2}{\left\langle \left| E_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) \right|^2 + \left| E_{\text{rd-b}}(\mathbf{r}_c; t_0) \right|^2 \right\rangle_{\mathbf{r}_c}} + \\ 127 & \frac{\left\langle E_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) E_{\text{rd-ho}}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} \left\langle E_{\text{rd-b}}^*(\mathbf{r}_c; t_0) E_{\text{rd-b}}(\mathbf{r}_c; t_0 + \Delta t) \right\rangle_{\mathbf{r}_c}}{\left\langle \left| E_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) \right|^2 + \left| E_{\text{rd-b}}(\mathbf{r}_c; t_0) \right|^2 \right\rangle_{\mathbf{r}_c}^2} + \\ & \frac{\left\langle E_{\text{rd-ho}}^*(\mathbf{r}_c; \mathbf{r}_o) E_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} \left\langle E_{\text{rd-b}}(\mathbf{r}_c; t_0) E_{\text{rd-b}}^*(\mathbf{r}_c; t_0 + \Delta t) \right\rangle_{\mathbf{r}_c}}{\left\langle \left| E_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) \right|^2 + \left| E_{\text{rd-b}}(\mathbf{r}_c; t_0) \right|^2 \right\rangle_{\mathbf{r}_c}^2} \end{aligned} \quad (31)$$

128 It can be found from Eq.(31) that  $C_{\text{rd}}(\Delta\mathbf{p})$  contains both the correlation of noise and  
 129 of transmitted field of the hidden object. When the noise evolves rapidly with time, its  
 130 time correlation can be seen as a  $\delta$  function, therefore,

$$131 \quad \left\langle E_{\text{rd-b}}(\mathbf{r}_c; t_0) E_{\text{rd-b}}^*(\mathbf{r}_c; t_0 + \Delta t) \right\rangle_{\mathbf{r}_c} = \begin{cases} \left\langle I_{\text{rd-b}}(\mathbf{r}_c; t_0) \right\rangle_{\mathbf{r}_c} & , \quad \Delta t = 0 \\ 0 & , \quad \Delta t > 0 \end{cases} \quad (32)$$

132 Substituting (32) into (31),

$$133 C_{rd}(\Delta\boldsymbol{\rho}) = \begin{cases} 1 & , |\Delta\boldsymbol{\rho}| = 0 \\ \frac{\left| \left\langle E_{rd-ho}(\mathbf{r}_c; \mathbf{r}_o) E_{rd-ho}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\boldsymbol{\rho}) \right\rangle_{\mathbf{r}_c} \right|^2}{\left\langle |E_{rd-ho}(\mathbf{r}_c; \mathbf{r}_o)|^2 + |E_{rd-b}(\mathbf{r}_c; t_0)|^2 \right\rangle_{\mathbf{r}_c}^2} & , |\Delta\boldsymbol{\rho}| > 0 \end{cases} \quad (33)$$

134 It can be found from Eq.(33) that  $C_{rd}(\Delta\boldsymbol{\rho})$  is a piecewise function. When  $|\Delta\boldsymbol{\rho}|$   
135 increases gradually from 0, the noise causes  $C_{rd}(\Delta\boldsymbol{\rho})$  to drop sharply from 1 to  
136  $\left| \left\langle I_{rd-ho}(\mathbf{r}_c; \mathbf{r}_o) \right\rangle_{\mathbf{r}_c} / \left\langle I_{rd}(\mathbf{r}_c; \mathbf{r}_o) \right\rangle_{\mathbf{r}_c} \right|^2$ , it deviates significantly from  $C_{ho}(\Delta\boldsymbol{\rho})$ . In the  
137 Supplementary Figure 2,  $C_{rd}$  calculated from the experimental data is compared with  
138  $C_{ho}$ , which confirms the sudden drop behavior and shows the significant deviation  
139 between the two correlations.

140 In order to image hidden objects, the adverse effect of noise on  $C_{rd}(\Delta\boldsymbol{\rho})$  must be  
141 eliminated. We denote the cross-correlation of speckle intensity images with the noise  
142 elimination as  $\tilde{C}_{rd}(\Delta\boldsymbol{\rho})$ ,

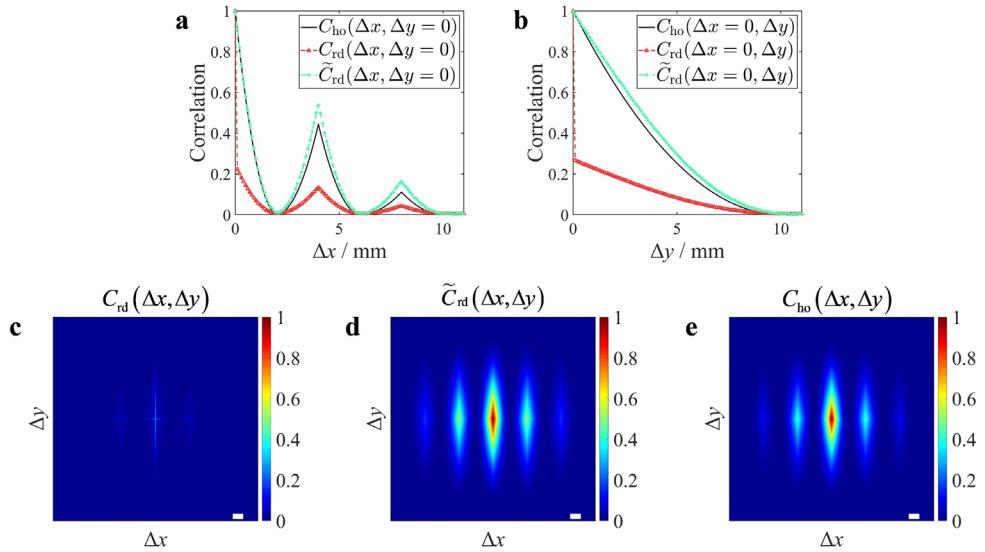
$$143 \tilde{C}_{rd}(\Delta\boldsymbol{\rho}) = \frac{\left| \left\langle E_{rd-ho}(\mathbf{r}_c; \mathbf{r}_o) E_{rd-ho}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\boldsymbol{\rho}) \right\rangle_{\mathbf{r}_c} \right|^2}{\left\langle I_{rd-ho}(\mathbf{r}_c; \mathbf{r}_o) \right\rangle_{\mathbf{r}_c}^2} \quad (34)$$

144 According to **Supplementary Note 1**,  $\tilde{C}_{rd}(\Delta\boldsymbol{\rho}) \approx C_{ho}(\Delta\boldsymbol{\rho})$ ,  $\tilde{C}_{rd}(\Delta\boldsymbol{\rho})$ , therefore, can  
145 be used to reconstruct the image of hidden objects. In order to recover  $\tilde{C}_{rd}(\Delta\boldsymbol{\rho})$ ,  
146  $\left| \left\langle E_{rd-ho}(\mathbf{r}_c; \mathbf{r}_o) E_{rd-ho}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\boldsymbol{\rho}) \right\rangle_{\mathbf{r}_c} \right|^2$  is extracted from  $C_{rd}(\Delta\boldsymbol{\rho})$ . By inverting Eq. (33),  
147 we have,

$$148 \left| \left\langle E_{rd-ho}(\mathbf{r}_c; \mathbf{r}_o) E_{rd-ho}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\boldsymbol{\rho}) \right\rangle_{\mathbf{r}_c} \right|^2 = \begin{cases} \left\langle I_{rd-ho}(\mathbf{r}_c; \mathbf{r}_o) \right\rangle_{\mathbf{r}_c}^2 & , |\Delta\boldsymbol{\rho}| = 0 \\ C_{rd}(\Delta\boldsymbol{\rho}) \cdot \left\langle I_{rd}(\mathbf{r}_c; \mathbf{r}_o) \right\rangle_{\mathbf{r}_c}^2 & , |\Delta\boldsymbol{\rho}| > 0 \end{cases} \quad (35)$$

149 From equation (35), it can be found that all values of

150  $\left| \left\langle E_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) E_{\text{rd-ho}}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} \right|^2$  are known except for  $\Delta\mathbf{p} = 0$ . Therefore,  
151  $\left\langle I_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) \right\rangle_{\mathbf{r}_c}^2$  can be extrapolated from the known values of  
152  $\left| \left\langle E_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) E_{\text{rd-ho}}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} \right|^2$ . Then  $\tilde{C}_{\text{rd}}(\Delta\mathbf{p})$  can be recovered with equation  
153 (34).



154  
155 Supplementary Figure 2: A comparison of  $C_{\text{ho}}$ ,  $C_{\text{rd}}$  and  $\tilde{C}_{\text{rd}}$ . Scale bar: 1mm.

156 We use quadratic polynomial to extrapolate  $\left\langle I_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) \right\rangle_{\mathbf{r}_c}^2$ . First, data points of  
157  $\left| \left\langle E_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) E_{\text{rd-ho}}^*(\mathbf{r}_c; \mathbf{r}_o + \Delta\mathbf{p}) \right\rangle_{\mathbf{r}_c} \right|^2$  in a monotonic decreasing region  
158  $(0 < \Delta x \leq l_x, 0 < \Delta y \leq l_y)$  near  $|\Delta\mathbf{p}|=0$  are chosen. Then these data points are  
159 substituted into the following quadratic polynomials,

$$160 \quad C_{\text{rd}}(\Delta x, \Delta y) \cdot \left\langle I_{\text{rd}}(\mathbf{r}_c; \mathbf{r}_o) \right\rangle_{\mathbf{r}_c}^2 = \left\langle I_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) \right\rangle_{\mathbf{r}_c}^2 + k_1(\Delta x + \Delta y) + k_2(\Delta x^2 + \Delta y^2) \quad (36)$$

161 Since the expression contains three unknown coefficients  $\left\langle I_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) \right\rangle_{\mathbf{r}_c}^2$ ,  $k_1$  and  $k_2$ ,  
162 at least three data points are required to solve these coefficients. We select three data  
163 points on x-axis, y-axis and 45° direction respectively, and substitute them into Eq.  
164 (36). Three groups of coefficients  $\left\langle I_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) \right\rangle_{\mathbf{r}_c}^2$ ,  $k_1$  and  $k_2$  are solved. Then three

165 values of  $\langle I_{\text{rd-ho}}(\mathbf{r}_c; \mathbf{r}_o) \rangle_{\mathbf{r}_c}^2$  are statistically averaged and substituted into Eq. (34) to  
166 obtain  $\tilde{C}_{\text{rd}}(\Delta\mathbf{p})$ .

167 The recovered  $\tilde{C}_{\text{rd}}$  is also drawn in Supplementary Figure 2. The good  
168 consistency between  $\tilde{C}_{\text{rd}}$  and  $C_{\text{ho}}$  indicates that the noise is successfully eliminated,  
169 and confirms the adaptability of the RSESC method in an environment with strong  
170 noise.

171

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